

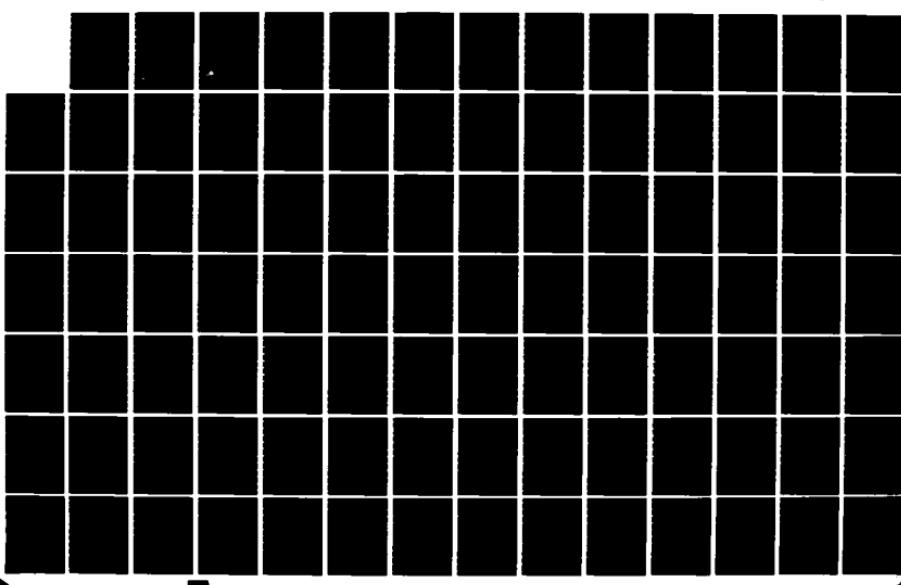
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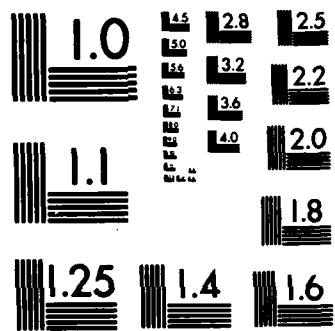
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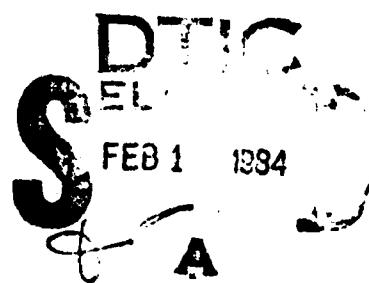
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FINAL REPORT 83-1155

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TROPICAL WEATHER SYSTEM AND OCEAN MODELING

FINAL REPORT SAI 83-1155

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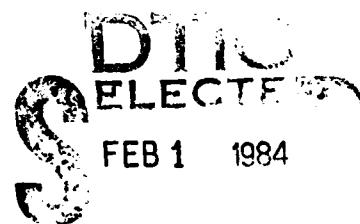
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Prepared Under:

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This final technical report documents the research efforts in the numerical simulations of tropical weather systems carried out under the auspices of ONR Contract N00014-82-C-2306, for the Atmospheric Physics Branch, Space Science Division, of the Naval Research Laboratory (NRL).

Our research efforts mainly concentrate in two major study areas: the interactions between two tropical cyclones and the construction of an axisymmetric ocean model. The research in both of these areas has been completed and reported in two articles. A paper entitled "A Numerical Study of the Interactions Between Two Tropical Cyclones" was accepted by the Monthly Weather Review and is to be published in the 1983 September issue. The draft of a second article on the ocean model has been completed. Both articles are included as Appendices in this final report. We have also attached a listing of the computer code of the ocean model, which is also stored in and accessible from the TI-ASC of NRL. Here, a brief summary of the study result will be given.

The interactions between atmospheric vortex pairs are simulated and studied with a nondivergent barotropic model and a three-dimensional tropical cyclone model (NRL/SAI mesoscale model). Numerical experiments with nondivergent barotropic vortex pairs show that the relative movements of the vortices are sensitive to the separation distance



and the characteristics of the swirling wind of the vortex. No observed mutual attraction is found in any of the nondivergent, barotropic vortex pairs tested.

Results from the 3D NRL/SAI tropical cyclone model show that on a constant-f plane with no mean wind, the movements of the two interacting tropical cyclones consist of a mutual cyclonic rotation, attraction, and eventual merging, in agreement with Fujiwhara's description. The displacement of one interacting storm in the mutual rotation is proportional to the combined strength of the binary system, but inversely proportional to the size of the storm and to the square of the separation distance. The rate of merging is related to the development of a mean secondary circulation on the radial-vertical plane, and is quite independent of the strength of the two tropical cyclones.

The latitudinal variation of the Coriolis parameter adds a northwest beta drift to the trajectories. Depending on their relative strength and location, the beta drift can either speed up the merging process or separate the two interacting tropical cyclones.

The axisymmetric ocean model consists of primitive equations for the conservation of momenta in three spatial dimensions and the buoyancy. A Boussineq assumption is made so that the background stratification is kept constant, the horizontal and vertical diffusion is of the Fickian type.

A leapfrog temporal integration is employed. The grid is fully staggered as Arakawa C type. The system is non-hydrostatic, the resultant elliptic equation for the pressure is solved by a stabilized error vector propagation technique. The basic equations, the finite differencing form, and boundary conditions are discussed in detail in the attached Appendix.

APPENDIX I
A NUMERICAL STUDY OF THE INTERACTIONS BETWEEN
TWO TROPICAL CYCLONES

A Numerical Study of the Interactions Between
Two Tropical Cyclones

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June 1983

ABSTRACT

The interactions between atmospheric vortex pairs are simulated and studied with a nondivergent barotropic model and a three-dimensional tropical cyclone model.

Numerical experiments with nondivergent barotropic vortex pairs show that the relative movements of the vortices are sensitive to the separation distance and the characteristics of the swirling wind of the vortex. No mutual attraction is found in any of the nondivergent, barotropic vortex pairs tested.

Results from the 3D tropical cyclone model show that on a constant-f plane with no mean wind, the movements of the two interacting tropical cyclones consist of a mutual cyclonic rotation, attraction, and eventual merging, in agreement with Fujiwhara's description. The displacement of one interacting storm in the mutual rotation is proportional to the combined strength of the binary system, but inversely proportional to the size of the storm and to the square of the separation distance. The rate of merging is related to the development of a mean secondary circulation on the radial-vertical plane, and is quite independent of the strength of the two tropical cyclones.

The latitudinal variation of the Coriolis parameter adds a northwest beta drift to the trajectories. Depending on their relative strength and location, the beta drift either speeds up the merging process or separates the two interacting tropical cyclones.

1. INTRODUCTION

When two tropical cyclones are present simultaneously in the same region, it is often observed that they rotate around each other with decreasing separation between them in the absence of large scale wind flow (Fig. 1). The phenomenon was made well-known by Fujiwhara (1921), and is therefore referred to as the Fujiwhara effect. By laboratory experiment and geophysical observation, Fujiwhara (1923, 1931) demonstrated that the relative motion of two counterclockwise vortices was a counterclockwise rotation. Haurwitz (1951) examined several tropical cyclone pairs by introducing the concept of center of mass around which the two tropical cyclones rotate about each other. By approximating the circulation around a tropical cyclone with that of a Rankine vortex, Haurwitz (1951) derived a relationship between the rotation rate and the sum of the total mass circulation of the two tropical cyclones. Many discrepancies were found when he applied the relationship to observations. Haurwitz attributed the discrepancies to the influence of large scale flow and lack of data, which led to deficiency in analyses.

Hoover (1961) studied binary tropical cyclones in both the Atlantic and Western Pacific Oceans. He found that the interaction between tropical cyclone pairs in the Western Pacific Ocean agrees with Fujiwhara's description while those pairs in the Atlantic Ocean rotated in an anticyclonic sense. He suggested that the different large scale atmospheric flow patterns in the two basins may have caused the binary systems to behave differently. The influence of the large scale flow was also noted by Liu and Wang (1966). They found that two interacting tropical cyclones in the Western Pacific are not always attracted to each other when there are strong shears in the environmental

flow. Recently, Dong and Neumann (1982)¹ found that storm pairs exhibiting behavior most in accordance with Fujiwhara's description were located in the Intertropical Convergence Zone where horizontal shears in large scale flow are negligible. They suggest that the effects of environmental flows be filtered before the real Fujiwhara effects can be determined. But to define and remove the large scale flows from observational data is difficult to accomplish.

Over the 35 year period 1946-1981, storm pairs known to have interactions averaged 1.5 annually in Western North Pacific and 0.33 annually in the Atlantic (Dong and Neumann, 1982). The presence of binary interacting has been noted to have contributed to forecast errors of tropical cyclone tracks (Brand, 1970; Jarrell et. al.; 1978; Neumann, 1982). Forecasting as well as analyzing a single tropical cyclone is often hindered by the paucity of observational data in the tropical cyclone basin (Neumann, 1982); the presence of two storms in close proximity can further compound the difficulties.

The purpose of our study is to investigate the interactions between two tropical cyclones by numerical simulations. Because the spatial resolutions of the models are better than the current observational network, and because numerical models can be controlled to produce "clean" results void of undesirable factors, analyzing realistic numerical simulations can sometimes result in a better isolation and understanding of the phenomenon than can be achieved from an observational approach. In this paper we will first determine the role of vorticity advection between the two vortices. For this purpose, a nondivergent, barotropic model is introduced to test two types of vortex pairs with different

¹Dong, K. and C. J. Neumann, 1982: On the relative motion of binary tropical cyclones. Regional Scientific Conference on Tropical Meteorology, Tsukuba Ibaraki, Japan, Oct. 1982.

swirling winds. These barotropic tests will be presented in Section 2. In Section 3, three-dimensional simulations of the interactions between two diabatically driven tropical cyclones on a constant-f plane and with variable f will be discussed. Our findings will be summarized in Section 4.

2. INTERACTIONS BETWEEN NONDIVERGENT, BAROTROPIC VORTEX PAIRS

In this section, we investigate the interactions between nondivergent barotropic vortex pairs. Through the interactions of such vortex pairs, we can determine the contribution of horizontal advection of vorticity, because in such a system advection is the only mechanism for interaction. A description of the nondivergent, barotropic model will be presented first, and the experimental design and the results will then be discussed.

a. Nondivergent Barotropic Model

The simple non-divergent, barotropic model can be described as

$$\frac{\partial}{\partial t} \nabla^2 \psi = -\underline{v}_\psi \cdot (\nabla^2 \psi + f), \text{ and} \quad (1)$$

$$\underline{v}_\psi = \hat{k} \times \nabla \psi \quad (2)$$

where f is the Coriolis parameter, ψ is the stream function, \underline{v}_ψ is the nondivergent wind, and \hat{k} is a vertically pointing unit vector. The boundary conditions for (1) and (2) are Neumann, i.e., $\nabla^2 \psi = 0$ at boundaries. The model has 51×51 grid points with a uniform horizontal resolution of 50 km. The relevant elliptic equation

$$\nabla^2 \psi = \zeta \quad (3)$$

where the relative vorticity is defined by

$$\zeta = \hat{\mathbf{r}} \cdot \nabla \times \underline{\psi}_g \quad (4)$$

is solved by a stabilized error vector propagation method (Madala, 1978).

b. Experimental Design

The major application of the nondivergent barotropic model is to determine the effects of separation distance and the radial distribution of tangential winds on the interaction of the two vortices. Two kinds of wind distributions were tested. The first kind (type A) of vortex is defined by its cyclonic swirl wind v_o as function of radius r from the vortex center

$$v_o = \begin{cases} Ar \left(1 - \sin \frac{\pi r}{r_o}\right), & 0 \leq r \leq r_o \\ 0, & \text{otherwise} \end{cases} \quad (5)$$

where constant $A = 4 \times 10^{-4} \text{ s}^{-1}$ and $r_o = 400 \text{ km}$. Equation (5) yields a maximum swirl of $\sim 26 \text{ m s}^{-1}$ at $r = 150 \text{ km}$ and a maximum vorticity

$\zeta \approx 7.2 \times 10^{-4} \text{ s}^{-1}$ at $r = 0$. We note that there is a cutoff of v_o at r_o .

The second kind of vortex (type B) is defined as

$$v_o = Br \exp \left(-\frac{r^2}{r_e^2}\right) \quad (6)$$

where the e-folding distance r_e is 150 km. By letting the constant $B = 3.6 \times 10^{-4} \text{ s}^{-1}$, (6) yields a vortex with similar strength as that described by (5) with maximum swirl of $\sim 29 \text{ m s}^{-1}$ and a maximum $\zeta \approx 7 \times 10^{-4} \text{ s}^{-1}$. Type B vortex differs from Type A in that there is no cutoff of swirl. Fig. 2 compares the radial distributions of relative vorticities described by (5) and (6).

Four initial separation distances (300, 400, 600, and 1000 km) have been tested for each type of vortex pairs. All intergrations with the barotropic model are performed with constant $f = 4.37 \times 10^{-5} \text{ s}^{-1}$.

c. Results

Fig. 3 shows the trajectories of storm pairs having two types of swirl wind at four separation distances. It is very clear from Fig. 3 that the smaller the separation distance the faster the mutual transport. For instance, at a separation distance of 1000 km, neither type A vortex pair (with swirl cutoff at $r = 300 \text{ km}$) nor type B vortex pair can induce mutual motion. But at a separation distance of 400 km, they move at a speed of $\sim 400 \text{ km day}^{-1}$.

It is also evident that the mutually-induced motions of type A and type B vortices are very different, in spite of the values of constants for A and B which were chosen to give vortices of similar strength. Furthermore, the trajectories of type A vortex pairs are more anticyclonic. This may be a result of the fact that type B vortices have positive vorticities at $r \leq 200 \text{ km}$ whereas the vorticities of type A vortices change sign at $r = 150 \text{ km}$ (Fig. 2). Only vortex pairs at small separation distances rotate in a cyclonic fashion because they interact with positive shears. The motion of vortices in our model can only be caused by the advection of vorticity, the shear in one vortex can very much determine the movement of the other. These results indicate that the mutual motion of two interacting nondivergent, barotropic vortex pair are quite sensitive to the characteristics of the swirl winds.

In all the experiments illustrated in Fig. 3 the storm pairs drift apart, there is no mutual attraction as observed in some interacting typhoons. This suggests that the observed mutual attraction in typhoon pairs may be due to

the divergence and/or convergence that is not included in the barotropic model. Indeed, complicated diabatic processes in tropical cyclone such as long wave radiation, surface boundary layer effects and moist convection generate convergent flow in lower troposphere and divergent flow in upper troposphere. The irrotational component of the vortex circulation may be responsible for the occurrence of the observed cyclonic rotation and mutual attraction.

3. INTERACTIONS BETWEEN TROPICAL CYCLONE PAIRS

We have seen that nondivergent pairs do not cause a mutual motion similar to the description of Fujiwhara. The observed Fujiwhara effects may be due to dynamics that can only be resolved by a more complete model. To see that, we will simulate the interactions between two diabatically driven tropical cyclones with a three-dimensional model.

a. Three-dimensional Tropical Cyclone Model

The baroclinic model is identical to the one in Chang and Madala (1980) and Chang (1982), except for parameterization of the latent heating. The governing equations are in surface-pressure-weighted flux form for conservation of momentum, temperature and water vapor. The normalized pressure $\sigma = p/p_s$ is the vertical coordinate, where p_s is the surface pressure. The system is assumed hydrostatic. The bulk boundary layer parameterization is based on a generalized similarity theory (Chang, 1981). The model has 51×51 horizontal grid points with seven sigma layers in the vertical. The horizontal resolution is 1° in both the latitudinal and longitudinal directions. The east-west boundaries are cyclic. The boundary conditions at the north and south boundaries are such that the second derivatives of thermodynamic variables

normal to the boundaries vanish. In addition, diffusion coefficients are increased near the north and south boundaries to damp numerical noise there.

Kuo's parameterization was used in Chang and Madala (1980) and Chang (1982), but a prescribed heating is applied in this model as done by Anthes (1971) in a axisymmetric model. The heating rate here is defined as

$$\dot{Q}(r, \sigma) = \begin{cases} Q_0 \frac{\pi r}{2R} \sin[\pi(\sigma - 0.1)], & \text{for } r \leq R, 0.1 \leq \sigma \leq 0.9 \\ 0, & \text{otherwise} \end{cases} \quad (7)$$

where r is the distance between a grid point and the low pressure center, and $R = 300$ km is the limit of the heating function. Two values of Q_0 , 100 K day^{-1} and 200 K day^{-1} , have been used in various numerical experiments to define the weak and strong tropical cyclones, respectively. The vertical and horizontal distributions of the heating pattern described in (7) are illustrated in Figs. 4 and 5. The vertical heating distribution is similar to that of the differences between temperatures in convective clouds (T_c) and the environment (T) in a mean hurricane season sounding for the Gulf of Mexico as computed by a one-dimensional cloud model (Anthes, 1977, Fig 4a). The horizontal heating distribution agrees with the mean rainfall rate inferred from satellite observation in a typhoon (Adler and Rogers, 1977), except for the observed smooth fall-off at $r \geq 300$ km. No effort is attempted to simulate the eye because of the model horizontal resolution.

The heating prescribed by (7) nevertheless generates realistic circulations for tropical cyclones. Figure 6 shows the radial distribution of the quasi-steady wind speeds at the sixth ($\bar{\sigma} = 0.85$) and seventh ($\bar{\sigma} = 0.965$) model layers after 24 h of heating with $Q_0 = 200 \text{ K day}^{-1}$. The wind speeds have a peak at $r = 1^{\circ}$ and decrease gradually outward without discontinuity

at $r = R = 300$ km.

We note however that by using the prescribed heating in (7) the effects of the interaction between the two cyclones on the scale of the cumulus convection cannot be adequately simulated. In reality the momentum field in each storm, which affects the cumulus convection, can be modified by the proximity of another storm. A change in the cumulus convection in each storm may alter the storm's intensities, which can in turn affect the interaction between the two tropical cyclones. But these feedbacks may be secondary and are only important when the separation of the two storms is small. As a preliminary study, the more economical, prescribed heating is used to investigate the first order effects in the interaction.

b. Experimental Design

The tropical cyclone pairs in all numerical experiments (Table 1) with the 3D model are dynamically initialized by a 24h stationary heating at two locations, i.e., by applying (7) at two fixed grid points for 24h. Dong and Neumann (1982) found that in real cases when the separation distances are less than 11 degree of latitude, cyclonic rotation predominates. Therefore, the two fixed grid points for the stationary heating are set ten degree longitude apart in all experiments to ensure the occurrence of interaction. After the dynamic initialization period, the heating patterns are allowed to follow the low pressure centers. In Exps. 1-3 we simulate the Fujiwhara effects in zero large scale winds on a constant-f plane for strong-strong (Exp. 1) weak-weak (Exp. 2) and strong-weak (Exp. 3) storm pairs. Exp. 1 and 3 are repeated in Exp. 5 and 6 on a real variation of f. There is only one single tropical cyclone in Exp. 4 to help isolate the effect of the beta-drift. Unlike on a

Table 1: List of three-dimensional Numerical Experiments

Exp.	q (K day^{-1})	Storm A	Storm B	$f(\text{s}^{-1})$	Characteristic
1	200	200		4.37×10^{-5}	strong-strong interaction
2	100	100		"	weak-weak interaction
3	100	100		"	weak-strong interaction
4	200	-		variable f	beta drift
5	200	200		"	strong-strong interaction
6	100	200		"	weak-strong interaction
7	200	100		"	strong-weak interaction

constant-f plane where geophysical orientation is not meaningful, the interactions with real f for a weak (west)-strong (east) pair and strong (west)-weak (east) pair are quite different, as we shall see later, thus Exp. 7 is conducted to study the latter situation.

c. Results on a Constant-f Plane

Fig. 7 shows the surface pressure field at 24, 48, 72, and 96h for Exp. 1. The southward displacement of storm A (west) and the northward displacement of storm B (east) at 24h indicate that their interaction has already caused the two storms to begin to rotate cyclonically in spite of the stationary heating. The merging of the two storms progresses with the merging of the outer isobars as observed (Fig. 1). By 96h, only the 996 mb isobars show two separate low pressure centers. The pressure at the center point of the model decreases by 10 mb while the approaching of the two initial low pressure centers between 24-96h can only account for a pressure decrease of 2 mb. This indicates that the mutual rotation and merging involve dynamics more complicated than merely advective processes.

Exp. 1-3 are integrated with a constant f. Therefore the orientation has little meaning and the results are independent of the absolute initial positions of the storms. Fig. 8 shows the trajectories of the storm centers in Exp. 1, in which two strong model tropical cyclones are of the same strength. The trajectories show that the two storms rotate about each other in a cyclonic fashion before the coalescence at 102h. The two trajectories are symmetric about the center of mass, which coincides with the center of the model domain. Superimposed on the symmetric rotation is a convergence of the two tropical cyclones. The distance between the two storms decreases from ~1024 km at 24h to ~612 km at 96h. The symmetry remains until 102h when the two heating

patterns overlap and one single large area of low pressure is formed.

Exp. 2 is identical to Exp. 1 except that the heating rate is reduced by one half. The cyclonic trajectories (Fig. 9) are still remarkably symmetric about the center of mass. Because of the weaker heating, two identifiable centers still exist at 120h when they are only \sim 100 km apart. We note again that at small separation distances feedbacks between cumulus convection and the storm pair's interaction have been masked by the prescribed heating in our model.

The speeds at which the two tropical cyclones in Exp. 1 and 2 rotate around and approach each other are shown in Fig. 10. The tangential velocity of the cyclonic rotation in Exp. 1 increases from \sim 3 m s $^{-1}$ at 24-36h to well over 6 m s $^{-1}$ after 72h as the separation between the two tropical cyclones becomes small. The rotation speeds in Exp. 2 are about 1 m s $^{-1}$ slower than those in Exp. 1. However, the rate of convergence seems quite independent of the combined strength as indicated by the radial velocities in Fig. 10. The faster rotations between stronger pairs are evident observationally at small separation distances (Dong and Neumann, 1982). At larger separation distances, this relationship is not clear because the observational data contains environmental influences.

Fig. 11 shows the trajectories of the two storm center in Exp. 3, in which the maximum heating rate is storm A is only half that of storm B. We see that the two storms still rotate about each other cyclonically and that they still move toward each other. However, the trajectories are asymmetric and the weaker storm A moves much faster than the stronger storm B in a way similar to that of a binary celestial system in which the two bodies have different masses. The "mass" of a vortex is perhaps best expressed as the product of its mean

angular velocity and the square of an effective radius $\bar{\omega}R^2$. We will discuss this further later. This type of interaction between storms of different intensities has been observed between Typhoons Flossie and Grace in 1950 (Liu and Wang, 1966). Instead of being stationary, the center of rotation moved within a small area defined by the lines connecting the two storm centers at different times as in Fig. 10.

Results in Exps. 1-3 demonstrate that on an f-plane with no large scale wind, interactions between two tropical cyclones cause the two storms to rotate cyclonically, to attract each other and to coalesce eventually.

In order to examine the momentum fields associated with the interaction we transform the model 500 mb wind fields in Exp. 1 onto a polar grid with respect to the center of the model domain. We now define the azimuthal mean velocity as

$$\bar{v} = (\bar{v}_r, \bar{v}_\theta, \bar{v}_z) = \frac{1}{2\pi} \int_0^{2\pi} (v_r, v_\theta, v_z) d\theta \quad (8)$$

where v_r, v_θ, v_z are radial, tangential and vertical velocities on the polar grid, and θ is the azimuthal angle. Fig. 12 shows the mean vertical (upper), tangential (middle), and radial (lower) velocities for Exp. 1 at 24, 48, 72, and 96h. It is interesting that the mean momentum fields relative to the center of domain shown here are similar to those in weak but intensifying tropical disturbances (Hawkins and Rubsam, 1968). For example, at 24h there is a maximum mean tangential velocity of $\sim 4 \text{ m s}^{-1}$ at $r \sim 600 \text{ km}$ and a minimum of $\sim -2 \text{ m s}^{-1}$ at $r \sim 400 \text{ km}$, reflecting the cyclonic wind fields about the two storm centers. The maximum tangential velocity gradually increases to $\sim 16 \text{ m s}^{-1}$ and moves toward the center to a radius $\sim 350 \text{ km}$ at 96h. The maximum inflow also

develops from 2 m s^{-1} at $r \sim 600 \text{ km}$ at 24h to $\sim 4 \text{ m s}^{-1}$ at $r \sim 350 \text{ km}$ at 96h. The evolution of the mean vertical velocity includes the increasing magnitude and contracting radii of maximum upward motion. The development of the mean circulation is accompanied with the pressure decrease of 10 mb at the domain midpoint from 24-96h, as discussed before.

The development of the azimuthal mean circulation can also be illustrated by comparing the kinetic energy (KE) of the mean velocity (KEM) and the KE of the eddy velocity (KEE), where

$$\text{KEM} = \iint_0^{1000 \text{ km}} (\bar{v}_r, \bar{v}_\theta) \cdot (\bar{v}_r, \bar{v}_\theta) r dr d\theta \quad (9)$$

$$\text{KEE} = \iint_0^{1000 \text{ km}} (v'_r, v'_\theta) \cdot (v'_r, v'_\theta) r dr d\theta \quad (10)$$

$$\text{and } (v'_r, v'_\theta) = (v_r, v_\theta) - (\bar{v}_r, \bar{v}_\theta) \quad (11)$$

As shown by Fig. 13, the KEE, which can mostly be attributed to the circulations around the two centers, reaches a quasi-steady state after 36h. Meanwhile the KEM, representing the strength of the mean circulation as depicted in Fig. 10 around the center of rotation, steadily increases until the coalescence of the two tropical cyclones. The ratio KEE/KEM decreases from ~ 3 at 24h to less than 1 at 96h.

These analyses suggest that a mean circulation relative to the center of rotation develops due to the interaction of two tropical cyclones. This mean

circulation includes tangential, radial and vertical components resembling those associated with tropical cyclones. It is therefore not surprising that the trajectories of the two interacting storms are similar to the trajectories in the hurricane boundary layer (e.g. Anthes 1982, Fig. 4.6). Compared with the nondivergent, barotropic experiments in Section 2, it seems that the diabatic heating in two storms plays a crucial role for the merging of the two storms.

In additional numerical experiments, the surface friction was suppressed to test the frictional effects in the interaction. Results from these experiments were nearly identical to those presented for Exps. 1-3. We also halved the coefficients for the internal dissipation. For the same given heating rates, the interactions are nearly the same except for faster rotation rates, because the model cyclones were stronger with less internal friction. Therefore, neither surface nor internal friction seem to be critical processes in the interactions.

d. Relation of Mutual Rotation Rate to Bulk Parameters of the System

The results presented so far indicate generally that the rate of the mutual cyclonic rotation depends on the strength of the binary system and the separation distance of the two interacting storms. Perhaps by relating some extent parameters through laws governing solid body rotation, a simple description of the numerical results is attainable.

We now consider the rotation of the binary storms is similar to that of a dumbbell. The equation of motion for rotation states that the torque τ acting on the binary system is equal to the product of the rotational inertia of the system I and the angular acceleration $\dot{\omega}$ with respect to the axis of rotation, i.e.

$$\tau = I\dot{\omega} \quad (12)$$

If we let R be an effective radius of the storm and H be the scale height, the

mass of one storm can be approximated by $\rho_0 \pi R^2 H$. Because the radius of the mutual rotation is about $L/2$, the rotational inertia

$$I \propto \rho_0 (R_A^2 + R_B^2) H L^2 \quad (13)$$

where ρ_0 is a reference density and subscript A and B are pertinent for storms A and B, respectively.

The torque is equal to the cross product of a force and radius of rotation. The force involved in the interaction can be approximated by the advection, then it can be scaled by $\rho_0 \bar{V}^2 / L$, where \bar{V} is the velocity of the mean circulation defined by (8). Because the mean circulation depends on the combined strength of the two interacting cyclones, therefore $\bar{V} \propto \bar{v}_A + \bar{v}_B$, where \bar{v}_A and \bar{v}_B are the mean wind speed within the effective radius R_A and R_B , respectively. Thus, the torque is proportional to

$$\tau \propto LH \rho_0 (\bar{R}_A^2 + \bar{R}_B^2) (\bar{v}_A + \bar{v}_B)^2 L^{-1} \quad (14)$$

We note that the torque contains the dimension of the kinetic energy of the two storms, which is ultimately related to the applied heating \dot{Q} in our model. Substituting (13) and (14) into (12) and dropping the over-bars, we get

$$(v_A + v_B)^2 \propto \dot{\omega} L^2 \sim f \omega L^2 \quad (15)$$

In above, f^{-1} is selected as the time scale, so that $\dot{\omega} \sim \omega/T \sim f\omega$. For our purposes of examining numerical results of a limited domain model away from the equator where f remain nearly a constant, the selection of f^{-1} as a time scale is justifiable.

$$\omega \propto \frac{(v_A + v_B)}{fL^2} \quad (16)$$

Relationship (16) states simply that the rate of the mutual rotation is proportional to the combined kinetic energy of the two interacting tropical cyclone and inversely proportional to the square of the separation distance. In addition, the displacement of one storm should be inversely proportional to its size, because the radius of rotation is inversely proportional to the mass $\pi p_0 R^2 H$.

In applying (16) to the numerical results, the mean wind speed within the radius of the gale force wind (17 m s^{-1}) is used as v_A and v_B . Excluding data when the separation distance is smaller than the sum of two radii of the gale force, rotation rates of binary system at every 6h were compared with $(v_A + v_B)^2 / fL^2$ for Exps 1,2 and 3 (Fig. 14). It is clear the (16) is a good description of the results Exp. 1-3. The rotation rates ω and quantities $(v_A + v_B)^2 / fL^2$ have a correlation coefficient of 0.81. Therefore, our numerical results can to some extent be represented by surprisingly simple relationship (16).

It should be noted, however, that (16) is arrived through several simplifying assumptions. These include approximating the mutual rotation of two vortex in the atmosphere by using solid body mechanics and excluding the merging from consideration. While (16) yields good correlation, it is only an approximation of the rotation component of the interaction.

e. The Effects of Variation of the Coriolis Parameter

Exps. 4-7 were carried out with variable Coriolis parameter, which can produce northwestward drifts of tropical cyclones in the northern hemisphere (Adem, 1956; Anthes and Hoke, 1975; Madala and Piacsek, 1975). The velocity of the drift depends on the latitude and the cyclone's circulation. To examine the free drift of a single tropical cyclone in our model, we carried out Exp. 4. As shown by trajectory C in Fig. 15, the model tropical cyclone has an initial northward movement, but changes toward the northwest after 36h, similar to the results of Anthes and Hoke (1975). The 0-72h mean drift velocity is 1.18 m s^{-1}

toward the west and 1.37 m s^{-1} toward the north. The center at 72h is $\sim 6^\circ$ to the west and $\sim 6.5^\circ$ to the north of the initial position.

The latitudinal variation of the Coriolis parameter has a pronounced effect on the trajectories of the two interacting tropical cyclones. The trajectories of the two tropical cyclones with equal strength in Exp. 5 are shown in Fig. 15. The two storms merge much faster than Exp. 1 due to the faster northwest drift of the storm located to the south. At 87h only one large low pressure center is identifiable, while in Exp. 1, two low pressure centers still existed at 96h (Fig. 7). Instead of rotating around the + point in Fig. 8 as in Exp. 1, storm A moves toward the southeast then quickly turns toward the northeast, while storm B rapidly moves northwestward and rotates cyclonically with respect to storm A. The two storms eventually merge into one at 87h, with storm B having traveled a much larger distance from its initial position than storm A. The relative trajectories of A and B with respect to trajectory C (Exp. 4) are computed. The resultant relative trajectories (not shown) are nearly the same as those in Exp. 1 (Fig. 3), indicating that the trajectories A and B in Fig. 15 is nearly a linear combination of the trajectories in Fig. 3 and the beta drift.

Exp. 6 is to be compared with Exp. 3, where storm A is weaker than storm B. The trajectories of the storm centers in Exp. 6 (Fig. 16) again appear very different from those in Exp. 3. The stronger storm B shows more noticeable northwest drift than in Exp. 3. The weaker storm A rotates cyclonically toward the southeast at a much reduced rate and with a smaller radius, apparently due to the counteracting beta drift.

Most interesting is Exp. 7, in which storm A is stronger than storm B. From 24 to 72h, the weaker storm B moves cyclonically relative to storm A.

In the meantime storm A moves slowly toward the southwest nearly perpendicular to and away from storm B. The trajectories take a strange turn after 72h because the two storms now are close to the boundaries and start to influence each other through the east-west boundaries because of the cyclic boundary conditions there.

The distinctively different behavior between Exps. 5-7 can be explained by examining schematically the vectors of forces upon each storm. We let the northwest drift be proportional to the storm's intensity and size (Rossby, 1948; Adem, 1956) and the force due to the interaction be proportional to the combined strength of the binary system but inversely proportional to the strength of individual storm as discussed in Section 3c. Figure 17 shows the vectors and the resultant directions of movements for storm A at 24h of Exps. 5-7. In Exp. 5 both the beta drift and interaction (both the rotation and convergence are counted for) are strong, the movement of the storm is mostly due south as evident in Fig. 15. In Exp. 6, the beta drift is weaker but the interaction is the strongest, the movement is nearly along the vector of the interaction. In Exp. 7, the beta effect is strong while the interaction is weak, results in a slow movement of the storm away from storm B.

4. SUMMARY

The interactions between two mesoscale cyclonic vortices in the absence of large scale winds have been investigated with a nondivergent, barotropic model and a three-dimensional model. Model results indicate that the interactions between a nondivergent barotropic vortex pair are very different from those observed between a tropical cyclone pair, and that our three-dimensional simulations agree with the observed Fujiwhara phenomenon.

Two types of vortex pairs with various initial separation distances have been tested with the nondivergent, barotropic model. No mutual attraction is found in any of the cases tested. The curvature of the mutually induced rotation depends on the radial profile of swirl winds (or vorticities) of each vortex, the speed of mutually induced motion is a function of separation — the closer the two vortices, the faster they move. This is quite understandable, because in such a nondivergent barotropic model the two vortices can only interact by advection. These numerical experiments suggest that the observed Fujiwhara phenomenon is caused by a more complex mechanism than just vorticity advection.

Our simulations with a three-dimensional model reproduce observed Fujiwhara effects. The trajectories of simulated strong-strong, weak-weak, and weak-strong tropical cyclone pairs on a constant-*f* plane all consist of cyclonic rotations and mutual attractions. The rotation rate between two strong tropical cyclones is generally faster than that between a weak pair. The rate of convergence of a weak pair is not slower than that between a strong pair.

Additional analyses show that as the tropical cyclone pair start to interact, there forms a mean circulation about the center of mass of the two storms as the pressure there decreases more than can be expected by simple advective merging. The development of the mean circulation, consisting of a cyclonic tangential flow and an inward radial flow, resembles the circulation in weak but intensifying tropical disturbances. The kinetic energy of this mean circulation grows by a factor of four in 72h in one experiment, while the kinetic energy of the circulations associated with individual tropical cyclone remains relatively unchanged. It suggests that the development of a mean circulation on the vertical-radial plane relative to the center of mass of the interacting storm pairs is crucial in generating the cyclonic mutual rotation and merging.

A simple analysis points out that the displacement of one tropical cyclone interacting with another is proportional to the combined strength of the

vortex pair and inversely proportional to its own size and to the square of the separation distance. Our model results fit this description well except for cases when interacting storms become highly asymmetric about their own centers.

The latitudinal variation of the Coriolis parameter (beta effect) has a large influence on the trajectories of the interacting storm pairs. The beta effect causes a northwest shift and a faster merging of the two tropical cyclones of equal strength. The trajectories of two interacting tropical cyclones of equal strength have a northwest drift superposed on the symmetrical trajectories found on the constant-f plane. Observation studies showed that typhoon pairs sometimes drifted away from each other if there were strong shears in large scale flow (Liu and Wang, 1966; Dong and Neumann, 1982). This study indicated that differential beta drifts can also cause the two interacting tropical cyclones of different strength to diverge when the one initially located to the west is stronger.

These findings should not be accepted without caution because of several limitations of the numerical model. The model domain is perhaps too small for two tropical cyclones. In addition, the horizontal resolution of $\frac{1}{2}^{\circ}$ is only marginal for resolving realistically the smaller scale dynamics near the center. Being a uniform grid model, without decreasing the horizontal resolution, the model domain cannot be expanded due to limited computing resources. The cyclic boundary conditions created problems (as evident in Exp. 7) when two storms may have interacted with each other through the east-west boundaries. Perhaps the most serious limitation of our simulation is the heating prescribed a prior in the three-dimensional simulations, which may have masked the interactions between two adjacent tropical cyclones on the scale of cumulus convections. However, the development of the mean circulation about the center of mass of the two tropical

cyclones occurs at very early stage of the interaction when the separation is still large. This suggests that the detailed characteristics of cumulus convection in individual storm may not be important in setting up the cyclic rotation and mutual attraction. The use of the prescribed heating was justifiable except at small separations where the divergent-convergent pattern in each storm may be modified due to the proximity of another one.

In future research, a parameterized convective heating should be utilized to investigate the abovementioned secondary effect of the cumulus convection. In addition, the parameterized heating may react to large scale winds in a nonlinear fashion. Therefore, the nonlinear effects of the large scale winds on the interactions of two tropical cyclones also ought to be studied. The question of what is the maximum separation distance for storm pair to interact is also left for future studies when numerical models of tropical cyclone cover a larger domain are constructed.

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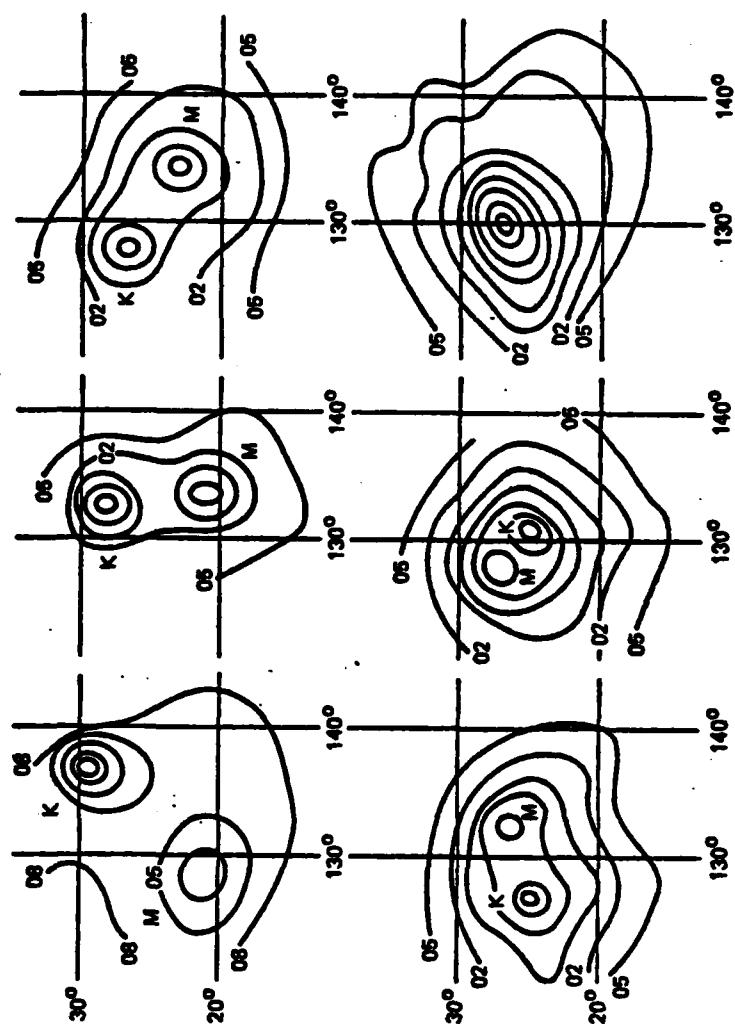
FIGURES

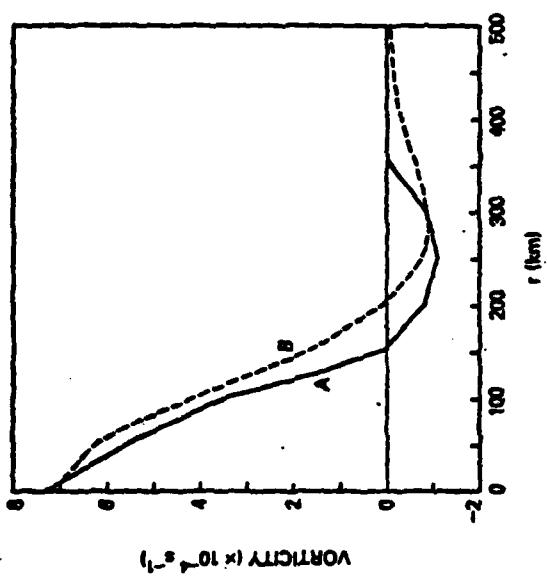
- Figure 1:** Surface isobaric analyses at 0000Z on Sept. 15-20, 1964 showing the rotation and merging of typhoons Kathy (K) and Marie (M). Isobars are plotted every 3 mb. (From Liu and Wang, 1966).
- Figure 2:** The radial distribution of relative vorticity for type A (solid) and type B vortex (dashed).
- Figure 3:** The trajectories of type A (solid lines) and type B (dashed lines) vortices at separation distances of 300, 400, 600, and 1000 km. The cross is the center of the domain. Time interval between two adjacent dots is 12h. Squares denote the initial vortex centers.
- Figure 4:** The vertical distribution of heating used in the model plotted on an arbitrary scale (solid), as compared with $(T_c - T)$ for mean hurricane sounding produced by a one-dimensional cloud model (Anthes, 1977, Fig. 4).
- Figure 5:** The horizontal distribution of heating used in the model plotted on an arbitrary scale (solid), as compared with the observed radial distribution of rainfall rate in a typhoon (Adler and Rodgers, 1977, Fig. 5).
- Figure 6:** The radial distribution of the quasi-steady wind speeds in model layer six ($\bar{\sigma} = 0.85$) and seven ($\bar{\sigma} = 0.963$) generated by the stationary prescribed heating.
- Figure 7:** The surface pressure field at 24, 48, 72, 84, and 96h. Contour intervals are 4 mb, the outmost closed isobars are 1008 mb. Longitudes are arbitrarily set.
- Figure 8:** The trajectories of storm centers in Exp. 1. Numbers on the curve denote times in hour.
- Figure 9:** As in Figure 8, except for Exp. 2.
- Figure 10:** The tangential and radial speeds of the two interacting tropical cyclones relative to their centers of mass in Exps. 1 and 2.
- Figure 11:** As in Fig. 8, except for Exp. 3.
- Figure 12:** The azimuthal mean radial (v_r), tangential (v_θ), and vertical velocity (w) of the wind fields relatives to the center of mass in Exp. 1 at 24, 48, 72 and 96h.
- Figure 13:** The development of the kinetic energy of the "mean" flow (solid) relative to the center of mass and the kinetic energy of the "eddy" associated with the two storm centers (dashed) in Exp. 1.
- Figure 14:** The rotational rates ω compared with $(v_A + v_B)^2/fL^2$.

Figure 15: The trajectories of the free drifting storm in Exp. 4 (Curve C) and of the two interacting storms in Exp. 5 (Curves a and B).

Figure 16: As in Fig. 15 except for Exps. 6 (solid lines with dots) and 7 (solid lines with squares).

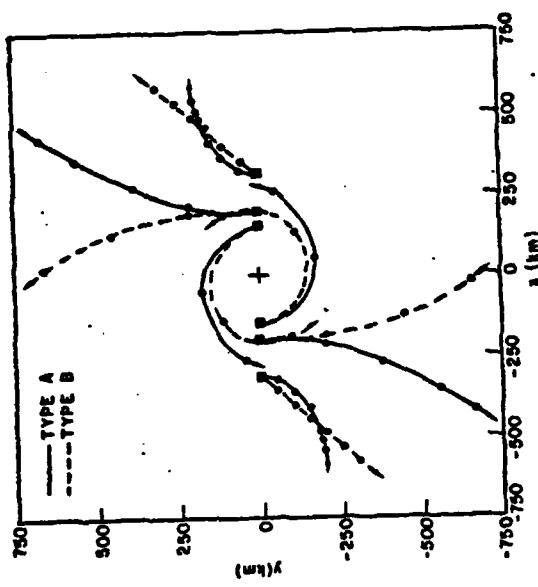
Figure 17: Vectors showing schematically the force of interaction (I), beta drift (D), and the resultant movement (M) for tropical cyclones A at 24h of Exps. 5, 6, and 7.

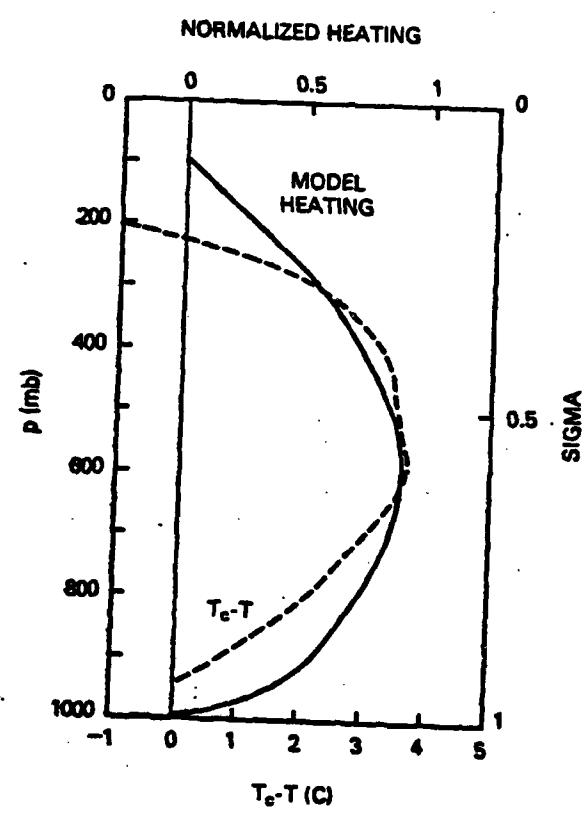


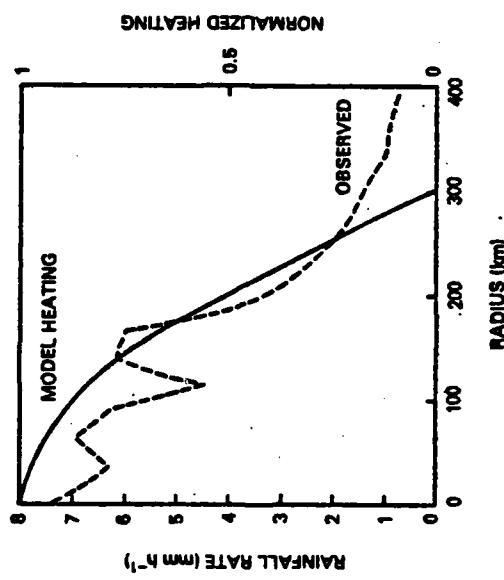


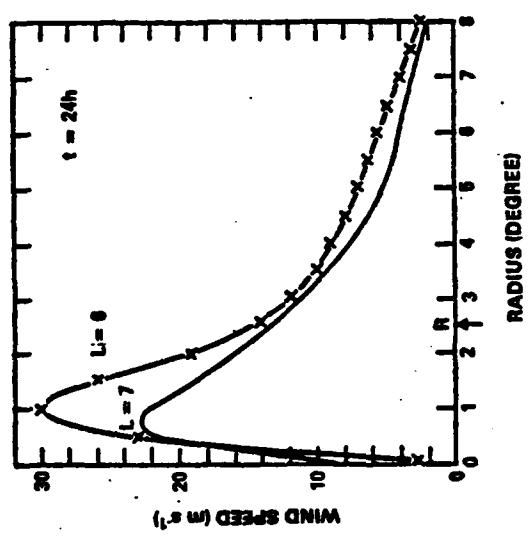
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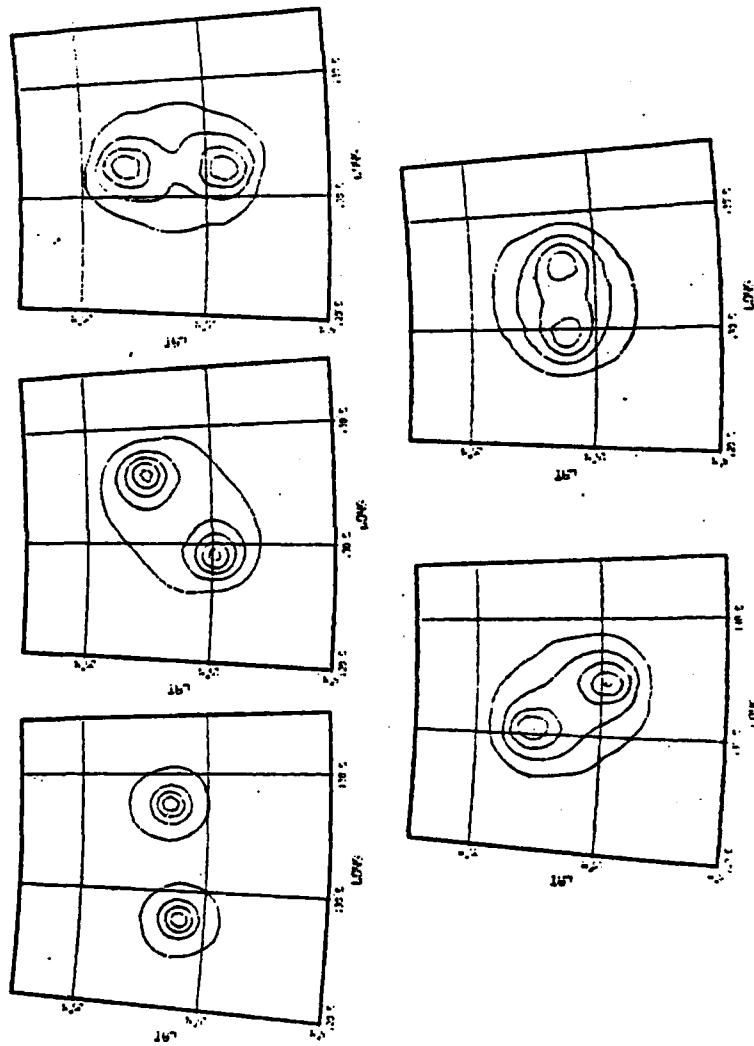


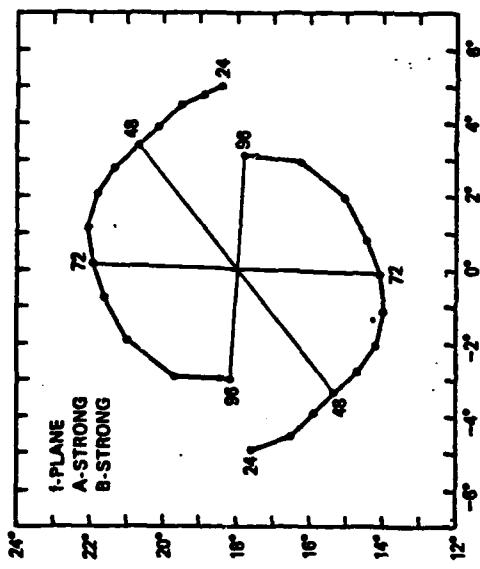




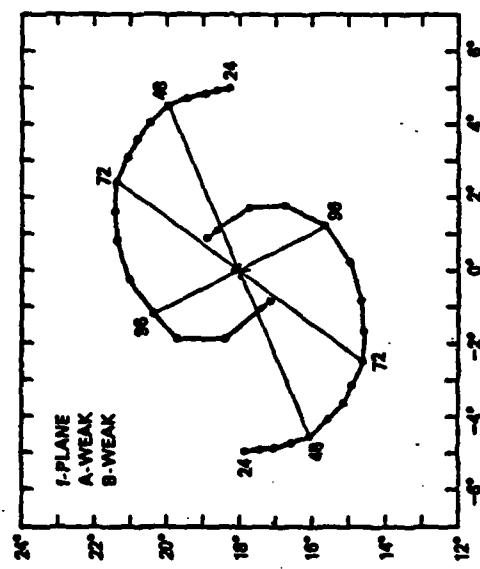


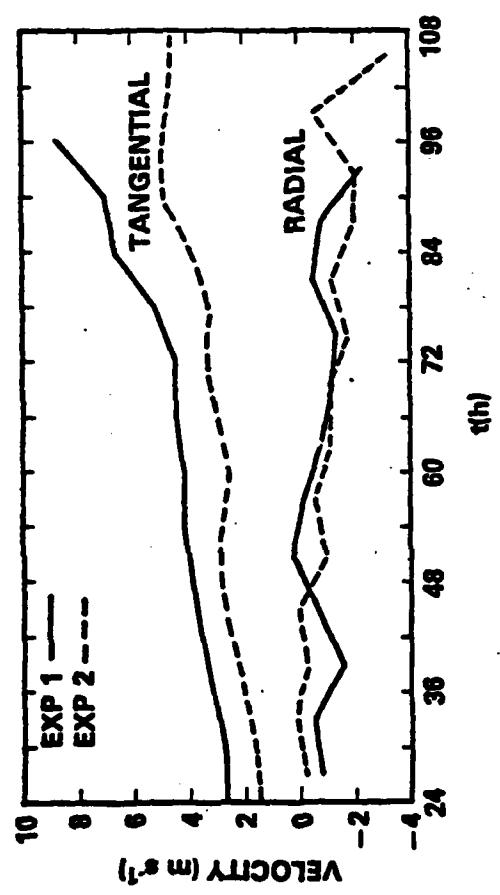
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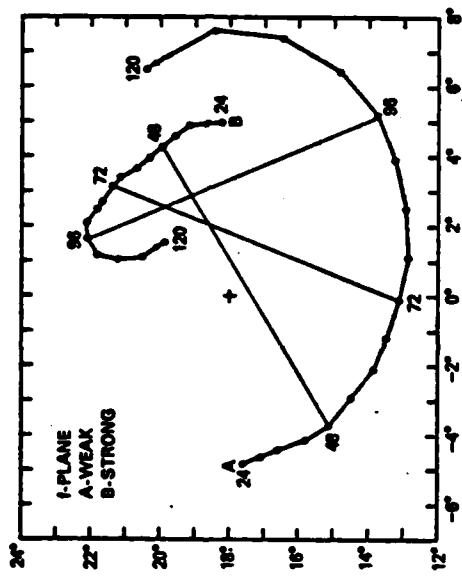


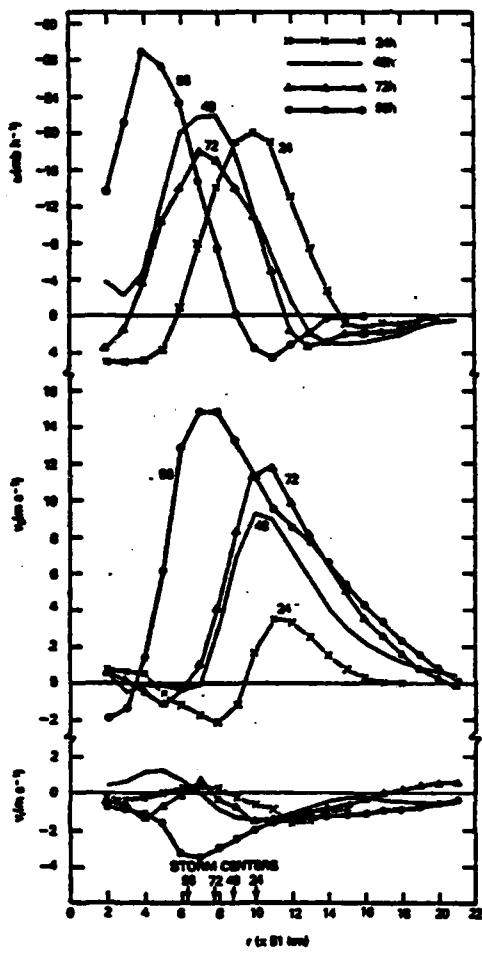
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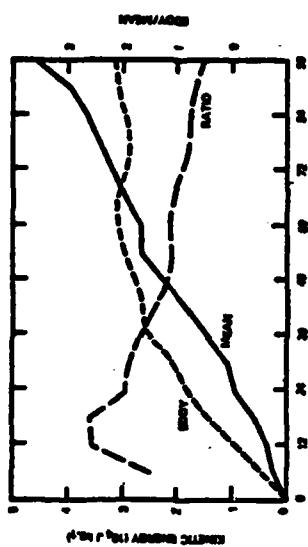


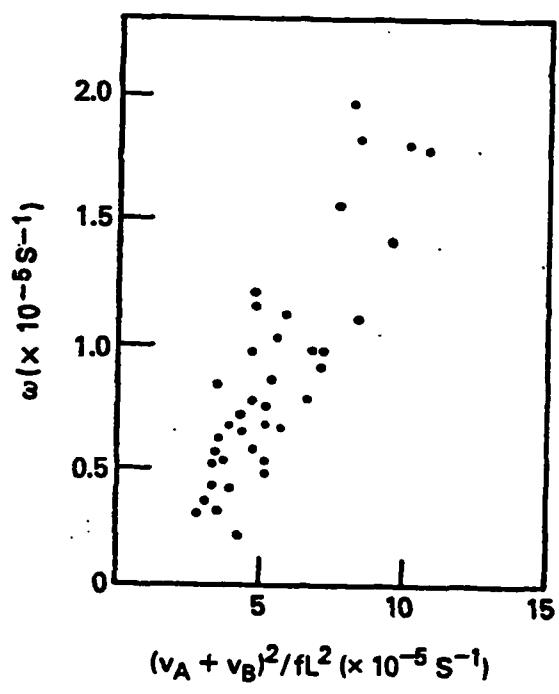
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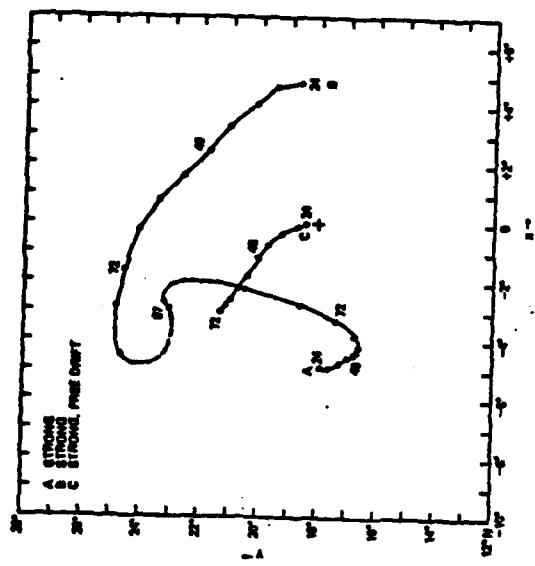


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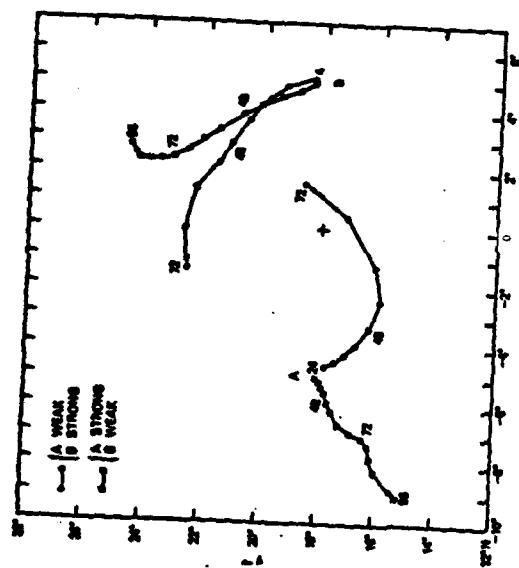




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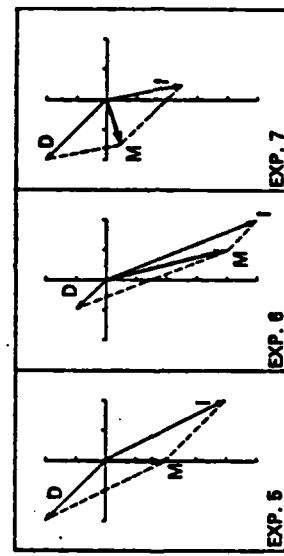


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APPENDIX II
AN AXISYMMETRIC MODEL FOR A NON-HYDROSTATIC
BOUSSINESQ OCEAN

AN AXISYMMETRIC MODEL FOR A NON-HYDROSTATIC
BOUSSINESQ OCEAN

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AN AXISYMMETRIC, NUMERICAL MODEL FOR A NON-HYDROSTATIC BOUSSINESQ OCEAN

1. GOVERNING EQUATIONS

The governing equations of the axisymmetric, non-hydrostatic, Boussinesq ocean model are

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial r} + w \frac{\partial u}{\partial z} = \frac{v^2}{r} + fv - \frac{1}{\rho_0} \frac{\partial p}{\partial r} + K_H \left(\nabla^2 u - \frac{u}{r^2} \right) + K_z \frac{\partial^2 u}{\partial z^2} \quad (1-1)$$

$$\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial r} + w \frac{\partial v}{\partial z} = - \frac{uv}{r} - fu + K_H \left(\nabla^2 v - \frac{v}{r^2} \right) + K_z \frac{\partial^2 v}{\partial z^2} \quad (1-2)$$

$$\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial r} + w \frac{\partial w}{\partial z} = - b - \frac{1}{\rho_0} \frac{\partial p}{\partial z} + K_H \nabla^2 w + K_z \frac{\partial^2 w}{\partial z^2} \quad (1-3)$$

$$\frac{\partial b}{\partial t} + u \frac{\partial b}{\partial r} + w \frac{\partial b}{\partial z} = N_z^2 w + K_H \nabla^2 b + K_z \frac{\partial^2 b}{\partial z^2} \quad (1-4)$$

where $\nabla^2 \equiv \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r}$, other symbols are listed in Appendix A.

Above, the density anomaly b is defined according to

$$b = \frac{\rho - \rho_r(z)}{\rho_0} g, \quad (1-5)$$

where $\rho_r(z)$ is a reference density and is a function of depth only. Brunt-Väisälä frequency N_z is defined as

$$N_z = \sqrt{\left(\frac{-g}{\rho_c} \frac{\partial \rho_r}{\partial z} \right)} \quad (1-6)$$

The continuity equation is that of the incompressible fluid,

$$\frac{1}{r} \frac{\partial u_r}{\partial r} - \frac{\partial w}{\partial z} = 0 \quad (1-7)$$

2. THE MODEL GRID

It is determined that a fully staggered grid is most expedient for storage economy for a given spatial resolution. As shown in Fig. 1, the radial (u) and the tangential (v) velocities are defined at cross points, vertical velocities (w) are defined at open circle points, and the pressures (p) and density anomalies (b) are defined at blackened dot points. This grid system has the following advantages:

- a) it saves storage for a given spatial resolution
- b) it is very economical in terms of number of computational operations for the finite difference (FD) equations of (1-1) to (1-4).
- c) it is very easy to specify the boundary conditions,
- d) the pressure diagnostic equation, of the elliptic type, can be reduced to the standard form, and
- e) there is no spatial separation of solutions on the grid.

In order to consistently index the grid points, we let index pair (ij) represent the i -th point in the r -direction and j -th point in the z -direction. In addition, m is the maximum number of points in the r -direction, and n , the maximum number of points in the z -direction. Therefore there are $m \times (n-1)$ points for radial and tangential velocities, $(m-1) \times n$ points for vertical velocities, and $(m-1) \times (n-1)$ points for mass distribution (b and p).

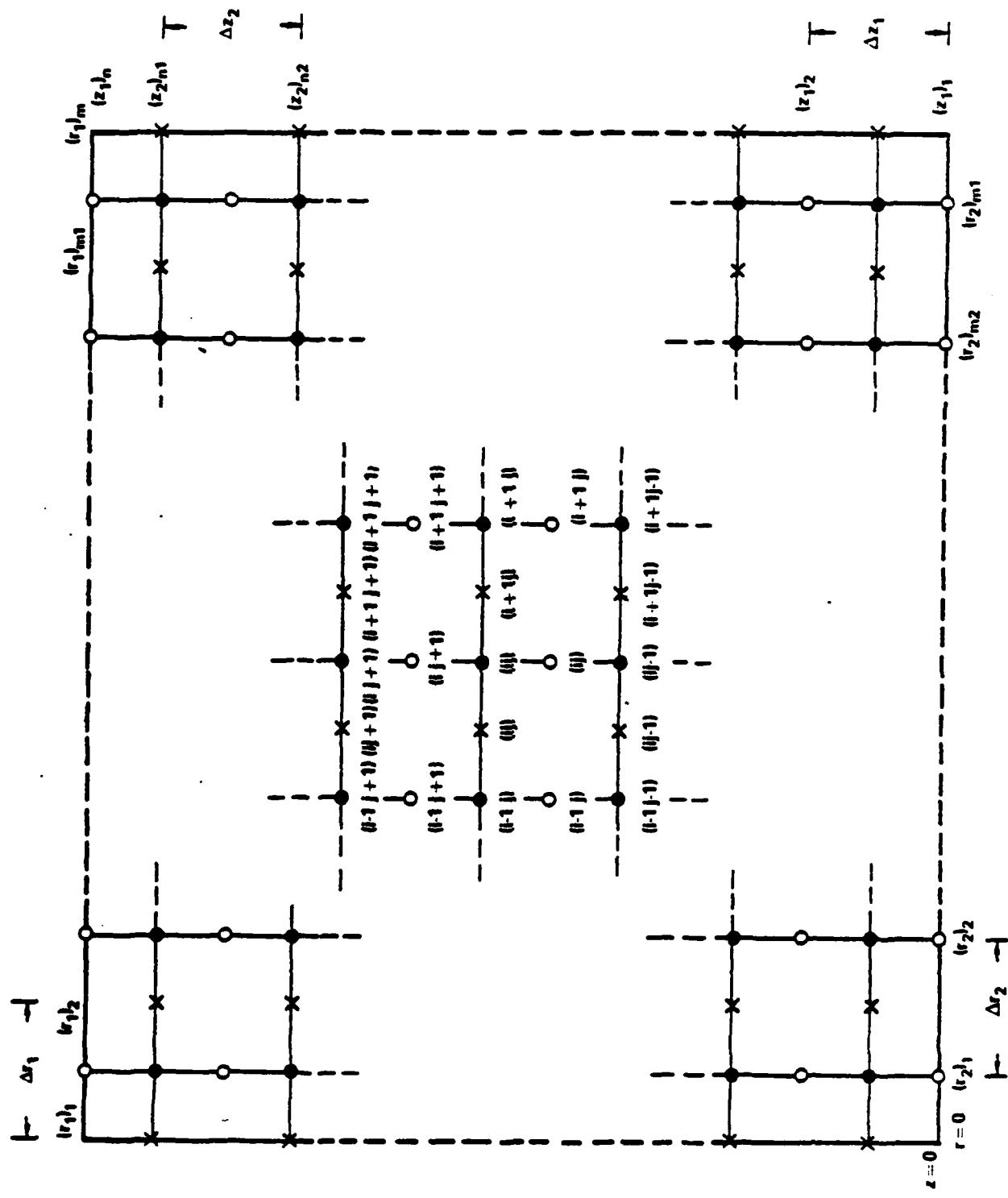


Fig. 1 The fully staggered grid system of the ocean model.

3. THE FINITE DIFFERENCE EQUATIONS

The leapfrog, or centered-in-time, integration scheme for the inviscid terms and the forward-in-time integration scheme for the viscous terms are used. The scheme is described as

$$\left\{ \begin{array}{l} u^{t+\Delta t} \\ v^{t+\Delta t} \\ w^{t+\Delta t} \\ b^{t+\Delta t} \end{array} \right\} = \left\{ \begin{array}{l} u^{t-\Delta t} \\ v^{t-\Delta t} \\ w^{t-\Delta t} \\ b^{t-\Delta t} \end{array} \right\} + 2\Delta t \left\{ \begin{array}{l} \frac{\partial u^t}{\partial t} \\ \frac{\partial v^t}{\partial t} \\ \frac{\partial w^t}{\partial t} \\ \frac{\partial b^t}{\partial t} \end{array} \right\} \quad (3-1)$$

A second order, or centered-in-space, scheme is applied to derive the tendencies in (3-1) according to (1-1) ~ (1-4).

(a) The Equation of Motion in r-direction

$$\frac{\partial u_{ij}^t}{\partial t} = H_{ij}^t - \frac{1}{\rho_0} \frac{1}{(\Delta r_2)_i} (p_{ij} - p_{i-ij}) \quad (3-2)$$

where

$$\begin{aligned}
H_{ij}^t = & -0.25 \left[\frac{1}{(\Delta r_1)_{i-1}} (u_{ij}^t + u_{i-1j}^t) (u_{ij}^t - u_{i-1j}^t) \right. \\
& + \frac{1}{(\Delta r_1)_i} (u_{i+1j}^t + u_{ij}^t) (u_{i+1j}^t - u_{ij}^t) \\
& + \frac{1}{(\Delta z_2)_j} (w_{i-1j}^t + w_{ij}^t) (u_{ij}^t - u_{ij-1}^t) \\
& \left. + \frac{1}{(\Delta z_2)_{j+1}} (w_{ij+1}^t + w_{i-1j+1}^t) (u_{ij+1}^t - u_{ij}^t) \right] \\
& + v_{ij}^t \left[\frac{v_{ij}^t}{(r_1)_i} + f \right] \\
& + K_H \left\{ \frac{1}{(\Delta r_2)_i} \left[\frac{1}{(\Delta r_1)_i} (u_{i+1j}^{t-\Delta t} - u_{ij}^{t-\Delta t}) \right. \right. \\
& \left. \left. - \frac{1}{(\Delta r_1)_{i-1}} (u_{ij}^{t-\Delta t} - u_{i-1j}^{t-\Delta t}) \right] \right. \\
& + 0.5 \left[\frac{1}{(r_2)_i (\Delta r_1)_i} (u_{i+1j}^{t-\Delta t} - u_{ij}^{t-\Delta t}) \right. \\
& \left. \left. + \frac{1}{(r_2)_{i-1} (\Delta r_1)_{i-1}} (u_{ij}^{t-\Delta t} - u_{i-1j}^{t-\Delta t}) \right] \right. \\
& \left. - \frac{u_{ij}^{t-\Delta t}}{(r_1)_i} \right\} + \frac{K_z}{(\Delta z_1)_j} \left[\frac{1}{(\Delta z_2)_{j-1}} (u_{ij+1}^{t-\Delta t} - u_{ij}^{t-\Delta t}) \right. \\
& \left. - \frac{1}{(\Delta z_2)_j} (u_{ij}^{t-\Delta t} - u_{ij-1}^{t-\Delta t}) \right]
\end{aligned}$$

3-3

(b) The Equation of Motion in ϵ -direction

$$\begin{aligned}
 \frac{\partial v_{ij}^t}{\partial t} = & -0.25 \left[\frac{1}{(\Delta r_1)_{i-1}} (u_{ij}^t + u_{i-1j}^t) (v_{ij}^t - v_{i-1j}^t) \right. \\
 & + \frac{1}{(\Delta r_1)_i} (u_{i+1j}^t + u_{ij}^t) (v_{i+1j}^t - v_{ij}^t) \\
 & + \frac{1}{(\Delta z_2)_j} (w_{i-1j}^t + w_{ij}^t) (v_{ij}^t - v_{ij-1}^t) \\
 & \left. + \frac{1}{(\Delta z_2)_{j+1}} (w_{ij+1}^t + w_{i-1j+1}^t) (v_{ij+1}^t - v_{ij}^t) \right] \\
 & - u_{ij}^t \left[\frac{v_{ij}^t}{(r_1)_i} + f \right] \\
 & + \frac{K_H}{(\Delta r_2)_i} \left\{ \left[\frac{1}{(\Delta r_1)_i} (v_{i+1j}^{t-\Delta t} - v_{ij}^{t-\Delta t}) \right. \right. \\
 & \left. \left. - \frac{1}{(\Delta r_1)_{i-1}} (v_{ij}^{t-\Delta t} - v_{i-1j}^{t-\Delta t}) \right] \right. \\
 & \left. + 0.5 \left[\frac{1}{(r_2)_i (\Delta r_1)_i} (v_{i+1j}^{t-\Delta t} - v_{ij}^{t-\Delta t}) \right. \right. \\
 & \left. \left. + \frac{1}{(r_2)_{i-1} (\Delta r_1)_{i-1}} (v_{ij}^{t-\Delta t} - v_{i-1j}^{t-\Delta t}) \right] \right. \\
 & \left. - \frac{v_{ij}^{t-\Delta t}}{(r_1)_i^2} \right\} - \frac{K_z}{(\Delta z_1)_j} \left[\frac{1}{(\Delta z_2)_{j+1}} (v_{ij+1}^{t-\Delta t} - v_{ij}^{t-\Delta t}) \right. \\
 & \left. - \frac{1}{(\Delta z_2)_j} (v_{ij}^{t-\Delta t} - v_{ij-1}^{t-\Delta t}) \right]
 \end{aligned}$$

5-4

(c)

The Equation of Motion in z-direction

$$\frac{\partial w_{ij}^t}{\partial t} = G_{ij}^t - \frac{1}{c_0(\Delta z_2)_j} (p_{ij} - p_{ij-1})$$

5-5

where

$$\begin{aligned}
 G_{ij}^t = & -0.25 \left[\frac{1}{(\Delta r_2)_i} (u_{ij}^t + u_{ij-1}^t) (w_{ij}^t - w_{ij-1}^t) \right. \\
 & + \frac{1}{(\Delta r_2)_{i+1}} (u_{i+1j}^t + u_{i+1j-1}^t) (w_{i+1j}^t - w_{ij}^t) \\
 & + \frac{1}{(\Delta z_1)_{j-1}} (w_{ij-1}^t + w_{ij}^t) (w_{ij}^t - w_{ij-1}^t) \\
 & \left. + \frac{1}{(\Delta z_1)_j} (w_{ij+1}^t + w_{ij}^t) (w_{ij+1}^t - w_{ij}^t) \right] \\
 & - 0.5 (b_{ij}^t + b_{ij-1}^t) \\
 & + K_H \left\{ \frac{1}{(\Delta r_1)_i} \left[\frac{1}{(\Delta r_2)_{i+1}} (w_{i+1j}^{t-\Delta t} - w_{ij}^{t-\Delta t}) \right. \right. \\
 & \left. - \frac{1}{(\Delta r_2)_i} (w_{ij}^{t-\Delta t} - w_{i-1j}^{t-\Delta t}) \right] \\
 & + 0.5 \left[\frac{1}{(r_1)_{i+1} (\Delta r_2)_{i+1}} (w_{i+1j}^{t-\Delta t} - w_{ij}^{t-\Delta t}) \right. \\
 & \left. \left. + \frac{1}{(r_1)_i (\Delta r_2)_i} (w_{ij}^{t-\Delta t} - w_{i-1j}^{t-\Delta t}) \right] \right\}
 \end{aligned}$$

$$+ \frac{K_z}{(\Delta z_2)_j} \left[\frac{1}{(\Delta z_1)_j} (w_{ij+1}^{t-\Delta t} - w_{ij}^{t-\Delta t}) - \frac{1}{(\Delta z_1)_{j-1}} (w_{ij}^{t-\Delta t} - w_{ij-1}^{t-\Delta t}) \right] \quad (S-6)$$

(d) The Thermodynamic Equation

$$\frac{\partial b_{ij}^t}{\partial t} = -0.5 \left[\frac{u_{ij}^t}{(\Delta r_2)_i} (b_{ij}^t - b_{i-1j}^t) + \frac{u_{i+1j}^t}{(\Delta r_2)_{i+1}} (b_{i+1j}^t - b_{ij}^t) \right]$$

$$+ \frac{w_{ij}^t}{(\Delta z_2)_j} (b_{ij}^t - b_{ij-1}^t) + \frac{w_{ij+1}}{(\Delta z_2)_{j+1}} (b_{ij+1}^t - b_{ij}^t) \Big]$$

$$+ 0.5 (w_{ij+1}^t + w_{ij}^t) N_z^2$$

$$+ K_H \left\{ \frac{1}{(\Delta r_1)_i} \left[\frac{1}{(\Delta r_2)_{i+1}} (b_{i+1j}^{t-\Delta t} - b_{ij}^{t-\Delta t}) - \frac{1}{(\Delta r_2)_i} (b_{ij}^{t-\Delta t} - b_{i-1j}^{t-\Delta t}) \right] \right.$$

$$+ 0.5 \left[\frac{1}{(\Delta r_1)_{i+1} (\Delta r_2)_{i+1}} (b_{i+1j}^{t-\Delta t} - b_{ij}^{t-\Delta t}) + \frac{1}{(\Delta r_1)_i (\Delta r_2)_i} (b_{ij}^{t-\Delta t} - b_{i-1j}^{t-\Delta t}) \right] \Big\}$$

$$+ \frac{K_z}{(\Delta z_1)_j} \left[\frac{1}{(\Delta z_2)_{j+1}} (b_{ij+1}^{t-\Delta t} - b_{ij}^{t-\Delta t}) - \frac{1}{(\Delta z_2)_i} (b_{ij}^{t-\Delta t} - b_{ij-1}^{t-\Delta t}) \right]$$

(S-7)

4. DERIVATION OF THE DIAGNOSTIC EQUATION FOR PRESSURE

The nonhydrostatic pressure at time t is needed to compute the pressure gradient forces in (3-2) and (3-5).

To "recover" the pressure from the motion fields, we make use of the continuity equation by differentiating (1-7) with time we get

$$\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial u}{\partial t} + \frac{\partial}{\partial z} \frac{\partial w}{\partial t} = 0 ,$$

which can be written in finite difference form for a mass point ij as

$$\frac{1}{\frac{1}{2}[(r_1)_i + (r_1)_{i+1}](\Delta r_1)_i} \left[(r_1)_{i+1} \frac{\partial u^t_{i+1j}}{\partial t} - (r_1)_i \frac{\partial u^t_{ij}}{\partial t} \right] \\ + \frac{1}{(\Delta z_1)_j} \left[\frac{\partial w^t_{ij+1}}{\partial t} - \frac{\partial w^t_{ij}}{\partial t} \right] = 0 \quad (4-1)$$

$$\text{Let } c_i = (r_1)_{i+1} / \left\{ \frac{1}{2} [(r_1)_i + (r_1)_{i+1}] (\Delta r_1)_i \right\},$$

(4-2)

$$\text{and } a_i = (r_1)_i / \left\{ \frac{1}{2} [(r_1)_i + (r_1)_{i+1}] (\Delta r_1)_i \right\}$$

Substituting (3-2), (3-5) and (4-2) into (4-1), we have

$$\begin{aligned}
c_i H_{i+1j}^t &= \frac{1}{\rho_0} \frac{c_i}{(\Delta r_2)_{i+1}} (p_{i+1j} - p_{ij}) + a_i H_{ij}^t \\
&+ \frac{1}{\rho_0} \frac{a_i}{(\Delta r_2)_i} (p_{ij} - p_{i-1j}) + \frac{1}{(\Delta z_1)_j} G_{ij+1}^t - \frac{1}{(\Delta z_1)_j} G_{ij}^t \\
&- \frac{1}{\rho_0 (\Delta z_1)_j (\Delta z_2)_{j+1}} (p_{ij+1} - p_{ij}) \\
&+ \frac{1}{\rho_0 (\Delta z_1)_j (\Delta z_2)_j} (p_{ij} - p_{ij-1}) = 0
\end{aligned}$$

After some rearrangements, we get

$$\begin{aligned}
&- \frac{c_i}{(\Delta r_2)_{i+1}} p_{i+1j} - \frac{a_i}{(\Delta r_2)_i} p_{i-1j} \\
&- \frac{1}{(\Delta z_1)_j (\Delta z_2)_j} p_{ij-1} - \frac{1}{(\Delta z_1)_j (\Delta z_2)_{j+1}} p_{ij+1} \\
&+ \left[\frac{c_i}{(\Delta r_2)_{i+1}} + \frac{a_i}{(\Delta r_2)_i} + \frac{1}{(\Delta z_1)_j (\Delta z_2)_{j+1}} + \frac{1}{(\Delta z_1)_j (\Delta z_2)_j} \right] p_{ij} \\
&= \rho_0 \left[-c_i H_{i+1j}^t + a_i H_{ij}^t - \frac{1}{(\Delta z_1)_j} G_{ij+1}^t + \frac{1}{(\Delta z_1)_j} G_{ij}^t \right] \quad (4-5)
\end{aligned}$$

Now let $F_{ij} = \text{RHS of } (4-5)$,

$$CX_i = \frac{c_i}{(\Delta r_2)_{i+1}} = (r_1)_i \cdot \left\{ \frac{1}{2} [(r_1)_i + (r_1)_{i+1}] \right. \\ \left. (\Delta r_1)_i (\Delta r_2)_{i+1} \right\} ,$$

$$AX_i = \frac{a_1}{(\Delta r_2)_i} = (r_1)_i \cdot \left\{ \frac{1}{2} [(r_1)_i + (r_1)_{i+1}] \right. \\ \left. (\Delta r_1)_i (\Delta r_2)_i \right\} ,$$

$$CZ_j = 1 / [(\Delta z_1)_j (\Delta z_2)_{j+1}] ,$$

$$AZ_j = 1 / [(\Delta z_1)_j (\Delta z_2)_j] , \text{ and}$$

$$BB_{ij} = - CX_i - AX_i - CZ_j - AZ_j \quad (4-4)$$

We obtain the standard form of an elliptic equation in FD form

$$AX_i p_{i-1,j} + AZ_j p_{ij-1} + BB_{ij} p_{ij} + CX_i p_{i+1,j} + CZ_j p_{ij+1} = F_{ij} \quad (4-5)$$

Equation (4-5) can be solved numerically by the SEVP solver (Madala, 1978), providing the boundary conditions are properly posed.

The conditions for the four boundaries are determined according to the following assumptions:

(a, At $(r_1)_i = (r_1)_1 = 0$, the natural condition for the cylindrical coordinates calls for $u = v = 0 = \partial u / \partial t = \partial v / \partial t$, the gradient balance requires that $(\partial P / \partial r)_{r=0} = 0$. Therefore an extra column of P is needed

$$P_{0,j} = P_{1,j}$$

(b) At $(r_1)_i = (r_1)_m$, assuming both the horizontal divergence and the vorticity are continuous, i.e., $\frac{\partial}{\partial r} \frac{1}{r} \frac{\partial u_r}{\partial r} = 0$ and $\frac{\partial}{\partial r} \frac{1}{r} \frac{\partial v_r}{\partial r} = 0$. These lead to

$$u_{mj} = b_a u_{m1j} + b_b [(r_1)_{m1} u_{m1j} - (r_1)_{m2} u_{m2j}] \quad (4-7)$$

$$v_{mj} = b_a v_{m1j} + b_b [(r_1)_{m1} v_{m1j} - (r_1)_{m2} v_{m2j}]$$

where $b_a = (r_1)_{m1}/(r_1)_m$, and

$$b_b = [(r_1)_{m1} + (r_1)_m] (\Delta r_1)_{m1} / \{(\Delta r_1)_{m2} [(r_1)_{m1} + (r_1)_{m2}] \}$$

Note that if b_b is set equal to zero, (4-7) describes a non-divergent and zero-vorticity boundary condition at $r = (r_1)_m$. Once v_{mj} is determined, a gradient balance at $r = (r_1)_m$ requires

$$\rho_0 v_{mj} \left[\frac{v_{mj}}{(r_1)_m} + f \right] = \frac{1}{(\Delta r_2)_m} (p_{mj} - p_{m-1j})$$

or

$$p_{mj} = p_{m-1j} + \rho_0 (\Delta r_2)_m v_{mj} \left[\frac{v_{mj}}{(r_1)_m} + f \right] \quad (4-8)$$

where a column of dummy points p_{mj} has been introduced for computational purposes. The second part of the RHS of (4-8) is thus the forcing function at $(r_1)_m$ for the elliptic equation (4-5).

(c) At the bottom, $w_{il} = \frac{\partial}{\partial t} w_{il} = 0$. Substituting these into the continuity equation (4-1), we get

$$\frac{1}{\frac{1}{2} (r_1)_i + (r_i)_{i+1} (\Delta r_1)_i} \left[(r_1)_{i+1} \frac{\partial u_{i+1}}{\partial t} - (r_1)_i \frac{\partial u_i}{\partial t} \right] \\ + \frac{1}{(\Delta z_1)_1} \frac{\partial w_{i2}}{\partial t} = 0 . \quad (4-9)$$

Following the same deduction between (4-1) and (4-5), we get an expression similar to (4-5) with the second term on the LHS and G_{il} in the RHS absent. Thus, P_{il} can be obtained by the same SEVP solver by setting $C_{il} = 0$ and $G_{il} = 0$.

(d) At top $w_{in} = \partial/\partial t w_{in} = 0$. Following the same line of reasoning as in (c), we obtain P_{in} by solving (4-5) with $C_{in} P_{in+1} = 0$ and $G_{in} = 0$.

In summary, the elliptic pressure diagnostic equation (4-4) is to be solved with the following boundary conditions

- 1) At $r = 0$ $P_{obj} = P_{lj}$ i.e., (4-6)
- 2) At $r = (r_1)_m$ $P_{mj} = P_{m-1j} + \text{function } (v_{mj})$ (4-8)
- 3) At $z = 0$ $A_{-1} = 0$ and $G_{il} = 0$
- 4) At $z = (z_1)_n$ $C_{in} = 0$ and $G_{in} = 0$

LIST OF SYMBOLS

Ax_i	an array of constants, varying only in r-direction, defined by (4-4), used in (4-5)
Az_j	an array of constants, varying only in z-direction defined by (4-4), used in (4-5)
a_i	an array of constants related to r_1 and Δr_1 used in (4-2)
Bb_{ij}	an array of constants, used in (4-5)
b	density anomalies, defined in (1-5), cm s^{-2}
Cx_i	an array of constants, varying only in r-direction, defined by (4-4), used in (4-5)
Cz_j	an array of constants, varying only in z-direction, defined by (4-4), used in (4-5)
c_i	an array of constants, related to r_1 and Δr_1 , used in (4-5)
f	Coriolis parameter, s^{-1}
g	gravitational acceleration, cm s^{-2}
i	an index, denoting i-th point in r-direction
j	an index, denoting j-th point in z-direction
K_H	horizontal diffusion coefficient, $\text{cm}^2 \text{s}^{-1}$
K_z	vertical diffusion coefficient, $\text{cm}^2 \text{s}^{-1}$
LHS	left hand side
m	the maximum number of grid points in r-direction, upper bound of i

m_1	$m-1$
m_2	$m-2$
N_z	Brunt-Väisälä frequency, s^{-1}
n	the maximum number of grid points in z -direction, upper bound of j
n_1	$n-1$
n_2	$n-2$
p	pressure, dyne cm^{-2}
RHS	right hand side
r	radius, cm
r_1	radii of momentum points, cm
r_2	radii of mass points, cm
Δr_1	distance between two horizontally adjacent momentum points, cm
Δr_2	distance between two horizontally adjacent mass points, cm
SEVP	<u>s</u> tabilized <u>e</u> rror <u>v</u> ector <u>p</u> ropagation
t	time, s
Δt	time interval, s
u	radial velocity, $cm s^{-1}$
v	tangential velocity, $cm s^{-1}$
w	vertical velocity, $cm s^{-1}$
z	height from ocean bottom, cm
z_1	heights of circle points, cm
z_2	heights of cross and dot points, cm

- Δz_1 distance between two vertically adjacent circle points, cm
- Δz_2 distance between two vertically adjacent cross or dot points
- ρ density, g cm^{-3}
- ρ_0 a constant density, 1 g cm^{-3}
- ρ_r a reference density, varying only in z -direction,
 g cm^{-3}

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REFERENCE

Madala, R. V., 1978: An Efficient Direct Solver for Separable and Non-Separable Elliptic Equations. Month Weather Review, 106, 1735-1741.

APPENDIX A - FORTRAN CODE FOR THE NON-HYDROSTATIC MODEL

A listing of FORTRAN code of the ocean model. The major functions of the main program and subroutines are as follows:

OCEAN	main program, calls all subroutines, manages job flow, controls input/output.
INIT	sets up independent variables, defines constants
START	defines initial conditions
{ PUTOUT	gets various fields ready for output
MAP	prints
ADVECT	computes all inviscous terms, except for the pressure gradient forces
DIFF	computes horizontal and vertical diffusions
PRESS	solves the pressure diagnostic equations and computes the pressure gradient forces, appears only in the non-hydrostatic version
MATINV	inverts matrices
BSM1	
BSM2	
BSM3	
FRWRD	matches forward
BOUNDV	sets outer boundary conditions for momentum
CHECK	checks if the time step is linearly stable.

} Used in SEVP method

APPENDIX B - FORTRAN CODE FOR THE HYDROSTATIC MODEL

The hydrostatic version of the model can be obtained by simplifying the non-hydrostatic version. In the hydrostatic version, the equation of motion in z -direction (1-5) is reduced to the hydrostatic equation

$$-\frac{1}{\rho_0} \frac{\partial p}{\partial z} = b \quad (B-1)$$

Instead of solving the elliptic equation (4-5), the pressure p is thus obtainable by vertically integrating (B-1). The vertical velocity w can also be computed by vertically integrating the continuity equation (1-7).

The FD forms of (B-1) and (1-7) are, respectively,

$$p_{ij} = p_{ij-1} - 0.5 \rho_0 (\Delta z_r)_j (B_{ij}^t + B_{ij-1}^t) \quad (B-2)$$

$$w_{ij} = w_{ij} + \frac{[(r_1)_{i+1} u_{i+1,j-1}^t - (r_1)_i u_{ij-1}^t]}{0.5 (\Delta r_1)_i [(r_1)_{i+1} + (r_1)_i]} (\Delta z_1)_j \quad (B-3)$$

<<SPLIT OCEAN,SOLRC0,PRINT,SEG

```

10      PROGRAM OCEAN
11      PARAMETER M21,M22
12      PARAMETER N10H=1,P2BH=2,N10H=3,N2BH=2
13      PARAMETER ND=2,PR1=M1+N0,M1=N1
14      DIMENSION DATA1(NC),DATA2(NC),DATA3(NC)
15      COMMON/ONE/VR1(M,N1),VT1(M,N1),VZ1(M,N1),B1(M,N1),VR2(M,N1),
16      VT2(M,N1),VZ2(M,N1),B2(M,N1),VR3(M,N1),VT3(M,N1),
17      VZ3(M,N1),B3(M,N1),P(M,N1)
18      EQUIVALENCE (DATA1,VR1),(DATA2,VR2),(DATA3,VR3)
19      DATA DATA1/ND=0./,DATA2/ND=0./,DATA3/ND=0./
20      COMMON/TMR/R1(M),R2(M),DR1(M),DR2(M),Z1(N),Z2(N),DZ1(N),DZ2(N) 0011000
21      COMMON/TMR/RMA,LPRC(M),BV2(N),ALPHA,BNDA,BNDB,CBRI,B,PK(N),ZK(N) 0012000
22      COMMON/PER/DELT,XTIME,ITIME,ISTEP,ISMC,ITAPE,TBV 0013000
23      CALL INDUMP 0014000
24      100 FORMAT(IJ) 0015000
25      READ(5,100)ITIME 0016000
26      READ(5,100)ITER 0017000
27      READ(5,100)INUT 0018000
28      READ(5,100)ISMC 0019000
29      ISSTEP=0 0020000
30      READ(5,100)ITAPE 0021000
31      CALL INIT 0022000
32      IF(ITIME.EQ.0)GP TO 10 0023000
33      C
34      C          CONTINUED INTEGRATION FROM A HISTORY TAPE 0024000
35      C
36      READ(1)ITIME,DATA1,DATA2,P 0025000
37      C
38      20 CALL START 0026000
39      20 XTIME=ITIME+3600. 0027000
40      C
41      C          PRINT OUT INITIAL FIELDS 0028000
42      C
43      CALL PUTOUT 0029000
44      IF(ITER.EQ.0)STEP 0030000
45      OR 90 ISSTEP=1,ITER 0031000
46      C
47      C          COMPLETE ALL INVISCID TERMS 0032000
48      C
49      CALL ADVECT 0033000
50      C
51      C          COMPUTE VISCOSITY TERMS 0034000
52      C
53      CALL DIFF 0035000
54      C
55      AND ACC PRESSURE GRADIENT FORCES TO TEMPERATURES 0036000
56      C
57      DIAGNOSE (RECOVER) THE PRESSURE FIELD 0037000
58      C
59      CALL PRESS 0038000
59      C
60      C          MARCHING IN TIME 0039000
61      C
62      FIRST TIME STEP IS FORWARD IF START IS CALLED 0040000
63      C
64      IF(ISTEP.EQ.1,AND,ITIME.EQ.0)DELT=0.5*DELT 0041000
65      CALL FORWARD 0042000
66      IF(ISTEP.EQ.1,AND,ITIME.EQ.0)DELT=2.*DELT 0043000
67      C
68      C          DEFINE BOUNDARY VALUES FOR VELOCITY 0044000
69      C
70      CALL BOUNDV 0045000
71      C
72      C          CHECK IF DELT IS STABLE 0046000
73      C
74      CALL CHECK 0047000
75      XTIME=XTIME+DELT 0048000
76      ITIME=ITIME+3600. 0049000
77      C
78      C          PRINT OUT RESULTS EVERY IBLT STEPS 0050000
79      C
80      IF(MOD(ISTEP,IBLT).EQ.0)CALL PUTOUT 0051000
81      C
82      C          WRITE HISTORY TAPE EVERY ITAPE STEPS 0052000
83      C
84      IF(MOD(ISTEP,ITAPE).EQ.0)WRITE(2)ITIME,DATA1,DATA2,P 0053000
85      C
86      CONTINUE 0054000
87      STOP 0055000
88      END 0056000
89      C
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1*      SUBROUTINE INIT
2*      PARAMETER NP2=1,NP2=1
3*      PARAMETER N1NP=1,N2NP=2,N1NP=1,N2NP=2
4*      COMMON/CNE/VB1(1,N1),VT1(1,N1),VZ1(1,N1),B1(N1,N1),VR2(1,N1),
5*      1          VT2(1,N1),VZ2(1,N1),B2(1,N1),VR3(1,N1),VT3(1,N1),
6*      2          VZ3(1,N1),B3(1,N1),B4(1,N1)
7*      COMMON/TB2/RH0,LH0R(N1),BV2(N1),ALPHA,BND4,BND5,C0P1,G,PW(N1),ZK(N1)
8*      COMMON/FOR/DELT,XTIME,TIME,ISTEP,ISNO,ITAPE,TBV
9*      PARAMETER NP3NP=1,NP3NP=1
10*     PARAMETER NBLKNP=2,NBLKNP=1,NBLKNP=1
11*     PARAMETER NP1NP=1,NP2NP=2,NP1NP=1,NP2NP=2
12*     REAL*8 RCOR,RINV,RINV1,RINV2,RINV3,RINV4,RINV5,RINV6,RINV7,RINV8,RINV9,
13*     COMMON/EVP/RINV(MP2,NP2,NBLK),RINV1(MP2,NP2,NBLK),RCOR(MP3),
14*     1          RTILT(MP2),F(MP,NP),NR8IZZ(NBLK),IS(NBLK),SLMP(NBLK),
15*     2          IE(NBLK),FI(MP),F14(MP),F21(NP),F2M(NP),AX(NP),AY(NP),
16*     3          PB(MP,NP),CX(MP),CY(MP)
17*     0001000
18* C      INITIALIZE ALL DEPENDENT VARIABLES AND CONSTANTS
19* C      0001000
20* C      0002000
21* C      0003000
22* C      ALPHA IS THE NONDIMENSIONAL SMOOTHING COEF,
23* C      FOR TIME SMOOTHING IN SUBROUTINE PRH0
24* C      0002000
25* C      0003000
26* C      DELTA=900,
27* C      ALPHAS=0.10
28* C      GS=80,
29* C      LAT=30
30* C      CR01=2.47,2722E-5=SIN(LAT*3.14159/180.)
31* C      DEFINE RADII AT GRID POINTS AND ALL GRID INTERVALS
32* C      00031000
33* C      00032000
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020 C
030 C      A AND B ARE CONSTANT USED IN SUBROUTINE BCLADV
040 C      FOR CONSTANT DIV. AND VERT. CONDITIONS
050 C
060 C      BNCA001(M1)/R1(P)
070 C      BNCA002(R1(M1)+R1(P))+DR1(M1)/((R1(M1)+R1(M2))+R1(M)+DR1(P2))
080 C      BNDB001,
090 C
100 C      DEFINE DENSITY RELATED CONSTANTS
110 C
120 C      RM001,
130 DD 130 J=1,N
140 130 BV2(J)=1,E-6
150 TBV001,
160 DD 135 J=1,N
170 135 TBV=MAX(TBV,BV2(J))
180 TBV=1/SQRT(TBV)
190 C
200 C      DEFINE HORIZONTAL AND VERTICAL DIFFUSION COEFFICIENTS
210 C
220 C      CREFM00,002*DR1(1)*2/DELT
230 C      COEFZ00,001*D21(1)*2/DELT
240 DD 140 I=1,M1
250 140 MK(I)=CREFM*(1.+5.*EXP(-FLOAT(M1-I)/7.))
260 DD 150 J=1,N1
270 150 ZK(J)=CREFZ*(1.+5.*EXP(-FLOAT(J-1)/5.)*EXP(-FLOAT(M1-J)/5.)))
280 RETURN
290 END

```

*** MEMBER START ***

```

10      SUBROUTINE START          0001000
20      PARAMETER N=21,N=21        0002000
30      PARAMETER M=10,M=10,M=10,M=20,M=20      0003000
40      COMMON/ONE/VR1(*,N),VT1(*,N),VZ1(*,N),P1(M1,N1),VR2(*,N1),
50      1      VT2(M,N1),VZ2(M1,N),B1(M1,N1),VR3(M,N1),VT3(M,N1),
60      2      VZ3(M1,N1),B2(M1,N1),VR4(M,N1),P(M1,N1)      0004000
70      COMMON/TW/R1(M,N),R2(M,N),DR1(M1),DR2(M1),Z1(N),Z2(N1),CZ1(N1),DZ2(N) 0005000
80      COMMON/TH/RH,RH,PRH(R1),VR2(N),ALPH4,BND4,CSP4,G,PK(N),ZK(N) 0006000
90      PARAMETER RH=2.0,M1=10,N=10,N1=10      0007000
100     DIMENSION DATA1(NE),DATA2(ND)          0008000
110     EQUIVALENCE (DATA1,VR1),(DATA2,VR2)          0009000
120   C          0010000
130   C          0011000
140   C          0012000
150   C          0013000
160   C          0014000
170   C          0015000
180   C          0016000
190   C          0017000
200   C          0018000
210   C          0019000
220   C          0020000
230   C          0021000
240   C          0022000
250   C          0023000
260   C          0024000
270   C          0025000
280   C          0026000
290   C          0027000
300   C          0028000
310   C          0029000
320   C          0030000
330   C          0031000
340   C          0032000
350   C          0033000
360   C          0034000
370   C          0035000
380   C          0036000
390   C          0037000
400   C          0038000
410   C          0039000
420   C          0040000
430   C          0041000
440   C          0042000
450   C          0043000
460   C          0044000
470   C          0045000
480   C          0046000
490   C          0047000
500   C          0048000
510   C          0049000
520   C          0050000
530   C          0051000
540   C          0052000
550   C          0053000
560   C          0054000
570   C          0055000
580   C          0056000
590   C          0057000
600   C          0058000
610   C          0059000
620   C          0060000
630   C          0061000
640   C          0062000
650   C          0063000
660   C          0064000
670   C          0065000
680   C          0066000
690   C          0067000
700   C          0068000
710   C          0069000
720   C          0070000
730   C          0071000
740   C          0072000
750   C          0073000
760   C          0074000
770   C          0075000
780   C          0076000
790   C          0077000
800   C          0078000
810   C          0079000
820   C          0080000
830   C          0081000
840   C          0082000
850   C          0083000
860   C          0084000
870   C          0085000
880   C          0086000
890   C          0087000
900   C          0088000
910   C          0089000
920   C          0090000
930   C          0091000
940   C          0092000
950   C          0093000
960   C          0094000
970   C          0095000
980   C          0096000
990   C          0097000
1000  C          0098000
1010  C          0099000
1020  C          0000000

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*** MEMBER BOUNDY

```

1*      SUBROUTINE BOUNDY          0001000
2*      PARAMETER M=2,N=1,N2=N*2   0002000
3*      PARAMETER M1=M-1,N2=N-2   0003000
4*      COMMON/VRL/VR1(M,N),VT1(M,N),VZ1(M,N),R1(M,N),VR2(M,N),   0004000
5*      VT2(M,N),VZ2(M,N),R2(M,N),VR3(M,N),VT3(M,N),   0005000
6*      VZ3(M,N),R3(M,N),P(M,N)   0006000
7*      COMMON/TRR/R1(M),R2(M),DR1(M),DR2(M),Z1(N),Z2(N),CZ1(N),DZ2(N) 0007000
8*      COMMON/TRR/RMA,LNCR(N),BV2(N),ALPHA,BNDA,BNDB,CORI,G,PK(N),ZK(N) 0008000
9* C                                     0009000
10* C      LATERAL BOUNDARY FOR TANGENTIAL AND RADIAL VELOCITIES 0010000
11* C      ASSUMING CONTINUOUS VORTICITY AND DIVERGENCE        0011000
12* C                                     0012000
13*      DO 10 J=1,N1               0013000
14*      VR2(M,J)=BNDA+VR2(M1,J)+BNCR*(R1(M1)*VR2(M1,J)+R1(M2)*VR2(M2,J)) 0014000
15*      10 VT2(M,J)=BNDA+VT2(M1,J)+BNDB*(R1(M1)*VT2(M1,J)+R1(M2)*VT2(M2,J)) 0015000
16*      RETURN                      0016000
17*      END                         0017000

```

*** MEMBER DIFF

```

1*      SUBROUTINE DIFF          0001000
2* C
3* C      COMPUTE THE DIFFUSION TERMS 0002000
4* C
5* C      PARAMETER M=21,N=21      0003000
6* C      PARAMETER M1=N-1,N2=N+2,N1=N-1,N2=N+2      0004000
7* C      CMMBN/PRM/VR1(M,N),VT1(M,N),VZ1(M1,N),B1(M1,N1),VR2(M,N1),      0005000
8* C      VT2(M,N1),VZ2(M1,N),B2(M1,N1),VR3(M,N1),VT3(M,N1),      0006000
9* C      VZ3(M1,N),B3(M1,N1),R(M1,N1)      0007000
10* C      CR4MBN/T=PR1(M),F2(M1),DR1(M1),DR2(M),Z1(N),Z2(N1),CZ1(N1),DZ2(N) 0008000
11* C      CR4MBN/THR/PRM,BPER(N1),BV2(N),ALPHA,BNDA,BNDB,CERI,G,HK(N),ZK(N) 0009000
12* C      DIMENSION VR(M,N1),VT(M,N1),VZ(M1,N),B(M1,N1)      0010000
13* C      EQUIVALENCE (VR,VR1),(VT,VT1),(VZ,VZ1),(B,B1)      0011000
14* C
15* C      HORIZONTAL DIFFUSION OF RADIAL VELOCITY      0012000
16* C
17* C      DD 10 JB1,N1      0013000
18* C      DD 10 IB2,M1      0014000
19* C      10 VR3(I,J)=VR3(I,J)+HK(I)*(((VR(I+1,J)-VR(I,J))/DR1(I))      0015000
20* C      *(VR(I,J)-VR(I-1,J))/DR1(I-1))/DR2(I)-VR(I,J)/(R1(I)*R1(I)) 0016000
21* C      1      *0.5*((VR(I+1,J)-VR(I,J))/DR1(I)*R2(I))      0017000
22* C      2      +(VR(I,J)-VR(I-1,J))/(DR1(I-1)*R2(I-1)))      0018000
23* C
24* C      HORIZONTAL DIFFUSION OF TANGENTIAL VELOCITY      0019000
25* C
26* C      DD 20 JB1,N1      0020000
27* C      DD 20 IB2,M1      0021000
28* C      20 VT3(I,J)=VT3(I,J)+HK(I)*(((VT(I+1,J)-VT(I,J))/DR1(I))      0022000
29* C      *(VT(I,J)-VT(I-1,J))/DR1(I-1))/DR2(I)-VT(I,J)/(R1(I)*R1(I)) 0023000
30* C      1      *0.5*((VT(I+1,J)-VT(I,J))/DR1(I)*R2(I))      0024000
31* C      2      +(VT(I,J)-VT(I-1,J))/(DR1(I-1)*R2(I-1)))      0025000
32* C
33* C      HORIZONTAL DIFFUSION OF VERTICAL VELOCITY      0026000
34* C
35* C      DD 30 JB2,N1      0027000
36* C      DD 30 IB2,M2      0028000
37* C      30 VZ3(I,J)=VZ3(I,J)+HK(I)*(((VZ(I+1,J)-VZ(I,J))/DR2(I+1)) 0029000
38* C      *(VZ(I,J)-VZ(I-1,J))/DR2(I))/DR1(I)      0030000
39* C      1      *0.5*((VZ(I+1,J)-VZ(I,J))/DR2(I+1)*R1(I+1))      0031000
40* C      2      +(VZ(I,J)-VZ(I-1,J))/(DR2(I)*R1(I)))      0032000
41* C      DD 40 JB2,N1      0033000
42* C      DD 40 IB2,M2      0034000
43* C      40 VZ3(I,J)=VZ3(I,J)+HK(I)*((VZ(2,J)-VZ(1,J))/(DR2(1)*DR1(1)) 0035000
44* C      1      *0.5*(VZ(2,J)-VZ(1,J))/(DR2(2)*R1(2)))      0036000
45* C      DD 50 JB2,N1      0037000
46* C      DD 50 IB2,M2      0038000
47* C      50 VZ3(M1,J)=VZ3(M1,J)+HK(M1)*((=VZ(M1,J)+VZ(M2,J))/(DR2(M1)*DR1(M1)) 0039000
48* C      1      +(=VZ(M1,J)+VZ(M2,J))/(DR2(M1)*R1(M1)))      0040000
49* C
50* C      HORIZONTAL DIFFUSION OF B      0041000
51* C
52* C      DD 60 JB1,N1      0042000
53* C      DD 60 IB2,M2      0043000
54* C      60 B3(I,J)=B3(I,J)+HK(I)*(((B(I+1,J)-B(I,J))/DR2(I+1)) 0044000
55* C      *(B(I,J)-B(I-1,J))/DR2(I))/DR1(I)      0045000
56* C      1      *0.5*((B(I+1,J)-B(I,J))/DR2(I+1)*R1(I+1))      0046000
57* C      2      +(B(I,J)-B(I-1,J))/(DR2(I)*R1(I)))      0047000
58* C      DD 70 JB1,N1      0048000
59* C      70 B3(I,J)=B3(I,J)+HK(I)*((B(2,J)-B(1,J))/(DR2(2)*DR1(1)) 0049000
60* C      1      *0.5*(B(2,J)-B(1,J))/(DR2(2)*R1(2)))      0050000
61* C      DD 80 JB1,N1      0051000
62* C      80 B3(M1,J)=B3(M1,J)+HK(M1)*((=B(M1,J)+B(M2,J))/(DR2(M1)*DR1(M1)) 0052000
63* C      1      +(=B(M1,J)+B(M2,J))/(DR2(M1)*R1(M1)))      0053000
64* C
65* C      VERTICAL DIFFUSION OF RADIAL VELOCITY      0054000
66* C
67* C      DD 90 JB2,N2      0055000
68* C      DD 90 IB2,M1      0056000
69* C      90 VR3(I,J)=VR3(I,J)+ZK(J)*((VR(I,J+1)-VR(I,J))/DZ2(J+1)) 0057000
70* C      1      -(VR(I,J)-VR(I,J-1))/DZ2(J))/DZ1(J)      0058000
71* C      DD 100 IB2,M1      0059000
72* C      100 VR3(I,J)=VR3(I,J)+ZK(1)*(VR(I,2)-VR(I,1))/(DZ2(2)*DZ1(1)) 0060000
73* C      110 VR3(I,N1)=VR3(I,N1)+ZK(1)*(=VR(I,N1)+VR(I,N2))/(DZ2(N1)*DZ1(N1)) 0061000
74* C
75* C      VERTICAL DIFFUSION OF TANGENTIAL VELOCITY      0062000

```

*** MEMBER DIFF

```

76*      DD 120 J=2,N2          0076000
77*      DD 120 I=2,M1          0077000
78*      120 VT3(I,J)=VT3(I,J)+ZK(J)*((VT(I,J+1)-VT(I,J))/DZ2(J+1)
79*           -(VT(I,J)-VT(I,J-1))/DZ2(J)/DZ1(J))          0078000
80*      DD 130 I=2,M1          0080000
81*      130 VT3(I,1)=VT3(I,1)+ZK(1)*((VT(I,2)-VT(I,1))/(DZ2(2)*DZ1(1))) 0081000
82*      DD 140 I=2,M1          0082000
83*      140 VT3(I,N1)=VT3(I,N1)+ZK(N1)*((VT(I,N1)-VT(I,N2))/(DZ2(N1)*DZ1(N1))) 0083000
84*      DD 150 I=2,M1          0084000
85* C
86* C          VERTICAL DIFFUSION OF VERTICAL VELOCITY        0085000
87* C
88*      DD 150 J=2,N1          0086000
89*      DD 150 I=M1,M1          0087000
90*      150 VZ3(I,J)=VZ3(I,J)+ZK(J)*((VZ(I,J+1)-VZ(I,J))/DZ1(J))
91*           -(VZ(I,J)-VZ(I,J-1))/DZ1(J-1)/DZ2(J))          0088000
92* C
93* C          VERTICAL DIFFUSION OF E                         0089000
94* C
95*      DD 160 J=2,N2          0090000
96*      DD 160 I=M1,M1          0091000
97*      160 E3(I,J)=E3(I,J)+ZK(J)*((E(I,J+1)-E(I,J))/DZ2(J+1)
98*           -(E(I,J)-E(I,J-1))/DZ2(J)/DZ1(J))          0092000
99*      DD 170 I=M1,M1          0093000
100*     170 H3(I,1)=H3(I,1)+ZK(1)*((H(I,2)-H(I,1))/(DZ2(2)*DZ1(1))) 0094000
101*     DD 180 I=M1,M1          0095000
102*     180 H3(I,N1)=H3(I,N1)+ZK(N1)*((H(I,N1)-H(I,N2))/(DZ2(N1)*DZ1(N1))) 0096000
103*     RETURN
104*     END

```

*** MEMBER FWDRC

```

1*      SUBROUTINE FWDRC
2*      PARAMETER M#21,N#21
3*      PARAMETER M1#1,M2#2,N1#1,N2#2
4*      PARAMETER NO#20,N1#1,M1#1,N1#1
5*      COMMON/RNE/DATA1(ND),DATA2(NC),DATA3(NC),P(M1,N1)
6*      COMMON/THR/RHO,THR(N1),BV2(N),ALPHA,BNDA,BNDB,CORI,G,MX(M),ZK(N)
7*      COMMON/FOR/DEL,TIME,ITIME,ISTEP,ISMC,ITARE,TBV
8* C
9* C          REPLACE DATA3 WITH THE NEW VALUES
10* C
11*      DD 10 IS1,ND
12*      10 DATA3(I)=DATA1(I)+2.*CELT*DATA3(I)
13* C
14* C          TIME SMOTHING
15* C
16*      IF(MAU(ISTEP,ISM0).NE.0)GO TO 30
17*      DD 20 IS1,ND
18*      20 DATA2(I)=DATA2(I)+(DATA1(I)+DATA3(I))/2.*DATA2(I))+ALPHA
19*      30 CONTINUE
20* C
21* C          FORWARD MARCHING
22* C
23*      DD 40 IS1,ND
24*      40 DATA1(I)=DATA2(I)
25*      DD 50 IS1,ND
26*      50 DATA2(I)=DATA3(I)
27* C
28* C          ZERO OUT DATA3 FOR NEXT STEP
29* C
30*      DD 60 IS1,ND
31*      60 DATA3(I)=0.
32*      RETURN
33*      END

```

SEE MEMBER CHECK

```

1*      SUBROUTINE CHECK
2*      PARAMETER N=21,N=21
3*      PARAMETER M=1,P=8M+2,N1=N=1,N2=N=2
4*      COMMON/ONE/V1(M,N),VT1(M,N),VZ1(M1,N),R1(M1,N1),VR2(M,N1),
5*      1          VT2(M,N1),VZ2(M1,N),R2(M1,N1),VR3(M,N1),VT3(M,N1),
6*      2          VZ3(M1,N),R3(M1,N1),P(M1,N1)
7*      COMMON/TMR/R1(M),R2(M1),DP1(M1),DR2(M),Z1(N),ZZ(N1),DZ1(N1),DZ2(N)
8*      COMMON/FOR/DELT,XTIME,ITIME,ISTEP,ISMN,ITAPE,TBV
9*      DIMENSION WORK1(N),WORK2(N)
10*     DATA 10 J=1,N1
11*     DATA 20 I=1,M
12*     20 WORK1(I)=DR2(I)/APAX1(1, VR2(I,J))
13*     MINBMINAG(WORK1)+1
14*     DT=WORK1(MIN)+0.9
15*     DT=AMINI(DT,DELT)
16*     10 CONTINUE
17*     DO 40 I=1,M1
18*     DO 30 J=1,N
19*     30 WORK2(J)=DZ2(J)/APAX1(1, VZ2(I,J))
20*     MINBMINAG(WORK2)+1
21*     DT=WORK2(MIN)+0.9
22*     DT=AMINI(DT,DELT)
23*     40 CONTINUE
24*     DT=AMINI(DT,TBV)
25*     IF(DT.GE.DELT)RETURN
26*     DELT=0.75*DELT
27*     PRINT 100,DELT
28* 100 FORMAT(////////,'*****DELT IS CHANGED TO',1PE11.2,'*****')
29*     RETURN
30*     END

```

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0002000
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0030000

*** MEMBER ZILCH

10	SUBROUTINE ZILCH(A,N)	0001000
20	DIMENSION A(N)	0002000
30	DO 10 I=1,N	0003000
40	10 A(I)=0.	0004000
50	RETURN	0005000
60	END	0006000

*** MEMBER ADVECT

```

14      SUBROUTINE ADVECT          0001000
24  C
34  C      COMPLETE THE ADVECTIVE TERMS 0002000
44  C
54  C      PARAMETER M21,M22           0003000
64  C      PARAMETER M10M1,M20M2,M10N1,M20N2 0004000
74  C      COMMON/BNE/VR1(M,N1),VT1(M,N1),VZ1(M1,N),B1(M1,N1),VR2(M,N1), 0005000
84  C      VT2(M,N1),VZ2(M1,N),B2(M1,N1),VR3(M,N1),VT3(M,N1), 0006000
94  C      1      VT2(M,N1),VZ2(M1,N),B2(M1,N1),VR3(M,N1),VT3(M,N1), 0007000
104 C      2      VZ3(M1,N),B3(M1,N1),P(M1,N1) 0008000
114 C      COMMON/TNO/R1(M),R2(M1),DR1(M1),DR2(M),Z1(N),Z2(N1),DZ1(N1),DZ2(N) 0009000
124 C      COMMON/THR/HR,GRHR(N1),BV2(N),ALPHA,BNDA,BNDB,C0(I,G,MK(M),ZK(N) 0010000
134 C      DIMENSION VR(M,N1),VT(M,N1),VZ(M1,N),B(M1,N1) 0011000
144 C      EQUIVALENCE (VR,VR2),(VT,VT2),(VZ,VZ2),(B,B2) 0012000
144 C
154 C      HORIZONTAL ADVECTION FOR RADIAL VELOCITY 0013000
164 C
174 C      DO 10 J=M1,N1 0014000
184 C      DO 10 I=M2,M1 0015000
194 C      10 VR3(I,J)=0.25*((VR(I,J)+VR(I-1,J))+(VR(I,J)+VR(I+1,J))/DR1(I-1) 0016000
204 C      1      +(VR(I+1,J)+VR(I,J))+(VR(I+1,J)+VR(I,J))/DR1(I)) 0017000
214 C      2      +VR3(I,J) 0018000
224 C
234 C      HORIZONTAL ADVECTION FOR TANGENTIAL VELOCITY 0019000
244 C
254 C      DO 20 J=M1,N1 0020000
264 C      DO 20 I=M2,M1 0021000
274 C      20 VT3(I,J)=0.25*((VR(I,J)+VR(I-1,J))+(VT(I,J)+VT(I-1,J))/DR1(I-1) 0022000
284 C      1      +(VR(I+1,J)+VR(I,J))+(VT(I+1,J)+VT(I,J))/DR1(I)) 0023000
294 C      2      +VT3(I,J) 0024000
304 C
314 C      HORIZONTAL ADVECTION FOR VERTICAL VELOCITY 0025000
324 C
334 C      DO 30 J=M2,N1 0026000
344 C      DO 30 I=M2,M2 0027000
354 C      30 VZ3(I,J)=0.25*((VR(I,J)+VR(I,J-1))+(VZ(I,J)+VZ(I-1,J))/DR2(I) 0028000
364 C      1      +(VR(I+1,J)+VR(I+1,J-1))+(VZ(I+1,J)+VZ(I,J))/DR2(I+1)) 0029000
374 C      2      +VZ3(I,J) 0030000
384 C      DO 40 J=M2,N1 0031000
394 C      40 VZ3(I,J)=0.25*(VR(2,J)+VR(2,J-1))+(VZ(2,J)+VZ(1,J))/DR2(2) 0032000
404 C      1      +VZ3(I,J) 0033000
414 C      DO 50 J=M2,N1 0034000
424 C      50 VZ3(M1,J)=0.25*(VR(M1,J)+VR(M1,J-1))+(VZ(M1,J)+VZ(M2,J))/DR2(M1) 0035000
434 C      1      +VZ3(M1,J) 0036000
444 C
454 C      HORIZONTAL ADVECTION FOR SUBVANCY 0037000
464 C
474 C      DO 60 J=M1,N1 0038000
484 C      DO 60 I=M2,M2 0039000
494 C      60 B3(I,J)=B3(I,J)+0.5*(VR(I,J)+B(I,J)+B(I-1,J))/DR2(I) 0040000
504 C      1      +VR(I+1,J)+B(I+1,J)+B(I,J))/DR2(I+1) 0041000
514 C      DO 70 J=M1,N1 0042000
524 C      70 B3(I,J)=B3(I,J)+0.5*VR(2,J)*(B(2,J)-B(1,J))/DR2(2) 0043000
534 C      DO 80 J=M1,N1 0044000
544 C      80 B3(M1,J)=B3(M1,J)+0.5*VR(M1,J)*(B(M1,J)-B(M2,J))/DR2(M1) 0045000
554 C
564 C      .      VERTICAL ADVECTION FOR RADIAL VELOCITY 0046000
574 C
584 C      DO 90 J=M2,N2 0047000
594 C      DO 90 I=M2,M1 0048000
604 C      90 VR3(I,J)+VR3(I,J)=0.25*((VZ(I-1,J)+VZ(I,J))+(VR(I,J)+VR(I,J-1)) 0049000
614 C      1      /DZ2(J)+(VZ(I,J+1)+VZ(I-1,J+1))+(VR(I,J+1)+VR(I,J))) 0050000
624 C      2      /DZ2(J+1)) 0051000
634 C      DO 95 I=M2,M1 0052000
644 C      95 VR3(I,I)+VR3(I,I)=0.25*(VZ(I,2)+VZ(I-1,2))+(VR(I,2)+VR(I,1)) 0053000
654 C      1      /DZ2(2) 0054000
664 C      DO 96 I=M2,M1 0055000
674 C      96 VR3(I,N1)=VR3(I,N1)=0.25*(VZ(I-1,N1)+VZ(I,N1))+(VR(I,N1)+VR(I,N2)) 0056000
684 C      1      /DZ2(N1) 0057000
694 C
704 C      VERTICAL ADVECTION FOR TANGENTIAL VELOCITY 0058000
714 C

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CARD IMAGE FILE ECITER(CIFER) -- VERSION 05.29 DATE 810/26/82 TIME 814:12:11 2194

*** MEMBER ADVECT

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72*      DA 100 J=2,*2          0072000
73*      DA 100 I=2,*1          0073000
74*      100 VT3(I,J)*VT3(I,J)=0.25*((VZ(I-1,J)+VZ(I,J))*(VT(I,J)+VT(I,J-1))
75*           /DZ2(J)+(VZ(I,J+1)+VZ(I-1,J+1))*(VT(I,J+1)+VT(I,J))) 0074000
76*           2 /DZ2(J+1)) 0075000
77*      DA 105 I=2,*1          0076000
78*      105 VT3(I,I)*VT3(I,I)=0.25*(VZ(I,2)+VZ(I=1,2))*(VT(I,2)+VT(I,1)) 0077000
79*           1 /DZ2(2) 0078000
80*      DA 106 I=2,*1          0079000
81*      106 VT3(I,N1)*VT3(I,N1)=0.25*(VZ(I,N1)+VZ(I=1,N1))*(VT(I,N1)+VT(I,N2)) 0080000
82*           1 /DZ2(N1) 0081000
83* C          VERTICAL ADVECTION FOR VERTICAL VELOCITY 0082000
84* C          0083000
85* C          0084000
86*      DA 140 J=2,*1          0085000
87*      DA 140 I=1,*1          0086000
88*      140 VZ3(I,J)*VZ3(I,J)=0.25*((VZ(I,J-1)+VZ(I,J))*(VZ(I,J)+VZ(I,J-1))
89*           1 /DZ1(J-1)+(VZ(I,J+1)+VZ(I,J))*(VZ(I,J+1)+VZ(I,J))) 0087000
90*           2 /DZ1(J)) 0088000
91* C          0089000
92* C          VERTICAL ADVECTION FOR E 0090000
93* C          0091000
94*      DA 150 J=2,*2          0092000
95*      DA 150 I=1,*1          0093000
96*      150 R3(I,J)*R3(I,J)=0.5*(VZ(I,J)*(B(I,J)+B(I,J-1))/DZ2(J)
97*           +VZ(I,J+1)*(B(I,J+1)+B(I,J))/DZ2(J+1)) 0094000
98*      DA 160 I=1,*1          0095000
99*      160 R3(I,I)*R3(I,I)=0.5*V2(I,2)*(R2(I,2)+R2(I,1))/DZ2(2) 0096000
100*     DA 170 I=1,*1          0097000
101*     170 R3(I,N1)*R3(I,N1)=0.5*VZ(I,N1)*(B(I,N1)+B(I,N2))/DZ2(N1) 0100000
102* C          INERTIA TERMS FOR HORIZONTAL MOMENTUM 0101000
103* C          0102000
104* C          0103000
105*      DA 110 J=1,N1          0104000
106*      DA 110 I=2,*1          0105000
107*      VR3(I,J)*VR3(I,J)+VT(I,J)*(VT(I,J)/R1(I)+CORI) 0106000
108*     110 VT3(I,J)*VT3(I,J)+VR(I,J)*(VT(I,J)/R1(I)+CORI) 0107000
109* C          0108000
110* C          BLODANCY TERM FOR VERTICAL ACCELERATION 0109000
111* C          0110000
112*      DA 120 J=2,*1          0111000
113*      DA 120 I=1,*1          0112000
114*     120 VZ3(I,J)*VZ3(I,J)=0.5*(R(I,J)+B(I,J-1)) 0113000
115* C          0114000
116* C          STRATIFICATION TERM 0115000
117* C          0116000
118*      DA 130 J=1,N1          0117000
119*      DA 130 I=1,*1          0118000
120*     130 R3(I,J)*R3(I,J)=0.5*(VZ(I,J)+BV2(J)+VZ(I,J+1)+BV2(J+1)) 0119000
121*     RETURN 0120000
122*     END 0121000
                                0122000

```

*** MEMBER ADVECT ADDED TO SOURCE -- 122 RECORDS

*** MEMBER PLTOUT

```

10      SUBROUTINE PLTOUT
11      PARAMETER M=21,N=21
12      PARAMETER M1=41,P2=84=2,N1=4=1,N2=84=2
13      COMMON/ONE/VR1(M,N),VT1(M,N),VZ1(M,N),B1(M,N),VR2(M,N),
14      1      VT2(M,N),VZ2(M,N),B2(M,N),VR3(M,N),VT3(M,N),
15      2      VZ3(M,N),B3(M,N),P(M,N)
16      COMMON/TIN/R1(M),R2(M),DR1(M),DR2(M),Z1(N),Z2(N),DZ1(N),DZ2(N)
17      COMMON/THR/RHM,PHMR(N),BV2(N),ALPHA,BNDR,CORI,G,MK(-),ZK(N)
18      COMMON/FOR/DELT,XTIME,ITIME,ISTEP,ISMC,ITAPE,TBV
19
20      C
21      C      THIS SUBROUTINE PRINT OUT FIELDS FOR A QUICK LOOK
22      C
23      DIMENSION IDUM(M,N)
24      700 FORMAT(//,,,' FACIAL VELOCITY (CM/S) AT T=',I6,' H')
25      705 FORMAT(//,,,' TANGENTIAL VELOCITY (CM/S) AT T=',I6,' H')
26      710 FORMAT(//,,,' VERTICAL VELOCITY (CM/S) AT T=',I6,' H')
27      715 FORMAT(//,,,' ELEVANCY FIELD (.,001) AT T=',I6,' H')
28      720 FORMAT(//,,,' PRESSURE (+10 DYN/CMS*2) AT TIME=',I6,' H')
29      725 FORMAT(1H1,/,20X,'***** OUTPUT AT TIME =',I6,' H',0H1,
30      1 FA,2,I' DAY ISTEP ',I7,' *****')
31      DAYTIME/86400,+L,0001
32      PRINT 720,ITIME,DAY,ISTEP
33      DD 10 J81,N1
34      DD 10 I81,M
35      10 IDUM(I,J)*VR2(I,J)
36      PRINT 700,ITIME
37      CALL MAP(IDUM,R1,22,M,N1)
38      DD 20 J81,N1
39      DD 20 I81,M
40      20 IDUM(I,J)*VT2(I,J)
41      PRINT 705,ITIME
42      CALL MAP(IDUM,R1,22,M,N1)
43      DD 30 J81,N
44      DD 30 I81,M1
45      30 IDUM(I,J)*VZ2(I,J)
46      PRINT 710,ITIME
47      CALL MAP(IDUM,R2,21,M1,N)
48      DD 40 J81,N1
49      DD 40 I81,M1
50      40 IDUM(I,J)*B2(I,J)+1.E3
51      PRINT 715,ITIME
52      CALL MAP(IDUM,R2,22,M1,N1)
53      DD 50 J81,N1
54      DD 50 I81,M1
55      50 IDUM(I,J)*P(I,J)+1.E-1
56      PRINT 720,ITIME
57      CALL MAP(IDUM,R2,22,M1,N1)
58      RETURN
59      END

```

*** REVERSE MAP

10	SUBROUTINE MAP(A,N,Z,MM,NN)	0001000
20	PARAMETER M21,M22	0002000
30	DIMENSION R(MM),Z(NN)	0003000
40	INTEGER A (M,N),IR(M),IZ(N)	0004000
50	70 FORMAT(1=5,7X,25I5)	0005000
60	80 FORMAT(1=5,I0,3X,25I5)	0006000
70	MM=M40(25,"")	0007000
80	DP 10 I81,MP	0008000
90	10 IR(I)=R(I)+1,E=5+R,1	0009000
100	DP 20 J81,NN	0010000
110	20 IZ(J)=Z(J)+1,E=E+0,1	0011000
120	PRINT 70	0012000
130	PRINT 70,(IR(I),I=1,MP)	0013000
140	PRINT 70	0014000
150	DP 30 JJ=1,NN	0015000
160	JAN=N+1=JJ	0016000
170	30 PRINT 80,IZ(J),(A(I,J),I=1,MP)	0017000
180	RETURN	0018000
190	END	0019000

SEE FENMEN PRESS

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10      SUBROUTINE PRESS          0001000
20      C                           0002000
30      C     THIS SUBROUTINE SETS UP FORCING FUNCTIONS AND BOUNDARY 0003000
40      C     CONDITIONS FOR THE PRESSURE DIAGNOSTIC EQUATIONS FOR 50004000
50      C                           0005000
60      C     PARAMETER M21,N21           0006000
70      C     PARAMETER N1=1,N2=2,N1B1=1,N2B1=2                   0007000
80      C     PARAMETER NBLK2,NBLK1=NBLK1=1                      0008000
90      C     REAL*8 RCNP,RINV,FINV1,RTILDA,DUMMV1               0009000
100     C     CRMON/CNE/VR1(1,N1),VT1(N1,N1),VZ1(N1,N),B1(M1,N1),VR2(N,N1),    0010000
110     1     VT2(N,N1),VZ2(N1,N),B2(N1,N1),VR3(N,N1),VT3(N,N1),    0011000
120     2     VZ3(N1,N),B3(N1,N1),P(M1,N1)                  0012000
130     C     CRMON/TAN/R1(M1),R2(M1),CR1(M1),DR2(N),Z1(N),Z2(N1),CZ1(N1),DZ2(N) 0013000
140     C     CRMON/TAN/R1(M1),R2(M1),CR1(M1),DR2(N),Z1(N),Z2(N1),CZ1(N1),DZ2(N) 0014000
150     C                           0015000
160     C     NP=NP IS THE SIZE OF X AND F                      0016000
170     C     NP=1+2,NP=N1+NBLK+1                            0017000
180     C                           0018000
190     C     PARAMETER NP=NP+1,RP=NP+1                     0019000
200     C     PARAMETER NP1=NP+1,NP2=NP+2,NP1=NP+1,NP2=NP+2       0020000
210     C     CRMON/EVP/RINV(MP2,NP2,NBLK),RINV1(MP2,NP2,NBLK1),PCOR(NP,3),   0021000
220     1     RTILLA(MP2),F(MP,NP),NBSIZZ(NBLK),IS(NBLK),SLNF(NBLK),    0022000
230     2     IE(NBLK),F11(MP),F1N(MP),F21(NP),F2H(NP),AX(NP),AY(NP),    0023000
240     3     PB(MP,NP),CX(NP),CY(NP)                      0024000
250     DIMENSION DUMMV1(MP2,PP2),X(NP,NP)                  0025000
260     EQUIVALENCE(DUMMV1,RINV1(1,1,NBLK))                0026000
270     PARAMETER MPNP=PP=NP                         0027000
280     DATA DUMMV1/MPM=0,/,X/MPM=0,/, 0028000
290     DATA NCALL/0/                                     0029000
300     C     NBSIZZ REPRESENTS NUMBER OF INTERIOR GRID POINTS IN EACH BLOCK IN X-DIMENSION 0030000
310     C     N2 REPRESENTS NUMBER OF INTERIOR GRID POINTS IN Y-DIRECTION 0031000
320     C     NBLK REPRESENTS NUMBER OF BLOCKS IN X-DIRECTION 0032000
330     C     THE VARIABLES A11,A1H,A21,A2H TAKES THE VALUE 0 FOR DIRICHLET B.C. 0033000
340     C     AND 1 FOR NEUMANN B.C. AT THEIR RESPECTIVE BOUNDARIES A11 CORRESPONDS TO 0034000
350     C     J=1 TO JN A1H TO JN A21 TO I=1 A2H TO I=MP 0035000
360     C     BOUNDARY CONDITIONS ARE 0036000
370     C     X(I,1)=(-A11)*X(I,1)+A11*(X(I,2)-F11(I)) 0037000
380     C     X(I,NP)=(-A1H)*X(I,NP)+A1H*(X(I,NP+1)-F1N(I)) 0038000
390     C     X(1,J)=(-A21)*X(1,J)+A21*(X(2,J)-F21(J)) 0039000
400     C     X(NP,J)=(-A2H)*X(NP,J)+A2H*(X(NP+1,J)-F2H(J)) 0040000
410     C     NCALL=NCALL+1 0041000
420     C                           0042000
430     C     DEFINE THE FORCING FUNCTION OF THE ELLIPTIC EQUATION 0043000
440     C                           0044000
450     DO 10 J=1,NP2 0045000
460     DO 10 I=1,MP2 0046000
470     10 F(I+1,J+1)=(CX(I+1)+DR2(I+1)+VR3(I+1,J)+VZ3(I,J+1))/CZ1(J) 0047000
480     1     -AX(I+1)+DR2(I)+VR3(I,J)+VZ3(I,J)/CZ1(J)+RN 0048000
490     C                           0049000
500     C     SET UP AN INITIAL GUESS 0050000
510     C                           0051000

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*** MEMBER PRESS

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524      IF(NCALL.GT.1)GP=30
530      CA 20 J=1,NP2
540      DA 20 I=1,NP2
550      20 X(I+1,J+1)=P(I,J)
560      30 CONTINUE
570      CALL ZILCH(F11,NP)
580      CALL ZILCH(F1N,NP)
590      CALL ZILCH(F21,NP)
600      C
610      C           DEFINE THE FORCING AT BOUNDARY SO THAT THERE
620      C           IS GRADIENT BALANCE AT OUTER BOUNDARY
630      C
640      D9 605 J=2,NP1
650      F2M(J)=R1H*D2(P)+VT2(H,J+1)+(VT2(H,J+1)/R1(H)*CBR1)
660      605 CONTINUE
670      A1NP1
680      A1NP1
690      A2NP1
700      A2NP1
710      DA 10P J=2,NP1
720      H2(2,J)=R2(2,J)+A2(2)*A21
730      F(2,J)=F(2,J)+A2(2)*F21(J)*A21
740      X(1,J)=X(1,0)*A21+X(1,J)
750      R=(NP-1,J)*R(NP-1,J)+C(X(NP-1))+A2H
760      F(NP-1,J)=F(NP-1,J)+C(X(NP-1))*F2H(J)*A2H
770      X(NP,J)=X(1,0)*A2H+X(NP,J)
780      101 CONTINUE
790      DA 102 I=2,NP1
800      R1(I,2)=R1H(I,2)+A1V(2)*A11
810      F(I,2)=F(I,2)+A1V(2)*F11(I)*A11
820      X(I,1)=X(1,0)*A11+A1(I)
830      R1(I,NP)=R1(I,NP-1)+C(Y(NP-1))+A1V
840      F(I,NP)=F(I,NP-1)-C(Y(NP-1))*F1N(I)*A1N
850      X(I,NP)=X(1,0)*A11+A1(I,NP)
860      102 CONTINUE
870      IF(NCALL.EQ.1)CALL BSH1
880      ENDROUTINE.BSH1,E=3
890      CALL BSH2(X,ERRER,A11,A1N,A21,A2H)
900      C
910      C           DEFINE THE DIAGNOSED PRESSURE
920      C
930      DA 110 J=1,N1
940      DA 110 I=1,N1
950      110 P(I,J)=X(I+1,J+1)
960      115 CONTINUE
970      C
980      C           ADD PRESSURE GRADIENT FORCES TO VR3 AND VZ3
990      C
1000     DA 120 J=1,N1
1010     DA 120 I=2,N1
1020     120 VR3(I,J)=VR3(I,J)+(P(I,J)-P(I-1,J))/(R1H*D2(P)(I))
1030     DA 130 J=2,N1
1040     DA 130 I=1,N1
1050     130 VZ3(I,J)=VZ3(I,J)+(P(I,J)-P(I,J-1))/(R1H*D2Z(J))
1060     RETURN
1070     END

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*** MEMBER B8P1

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14      SI BROUTINE R8M1          0001000
15      PARAMETER M821,M822          0002000
16      PARAMETER M8M=1,M8M=2,N18M=1,N28M=2          0003000
17      PARAMETER M8P4=1,M8P4=1          0004000
18      PARAMETER M818M=1,M828M=2,NP18NP=1,NP28NP=2          0005000
19      PARAMETER NBLK=M2,NBLK=M1          0006000
20      READ* FCSR,RIN1,RINV1,FTILDA,DUMMY1          0007000
21      COMMON/TMP/R1(M),R2(M1),DR1(M1),DR2(M),Z1(N),Z2(N1),DZ1(N1),DZ2(N) 0008000
22      COMMON/EVP/RINV(MP2,MP2,NBLK),RINV1(MP2,MP2,NBLK1),RCOR(MP,3), 0009000
23      RTILDA(MP2),F(MP,MP),NBSIZ2(NBLK),IS(NBLK),SLMP(NBLK), 0010000
24      1 IE(NELK),F1(MP),F1N(MP),F21(NP),F2M(NP),AX(MP),AY(NP), 0011000
25      2 0012000
26      3 0013000
27      0014000
28      0015000
29      0016000
30      0017000
31      0018000
32      0019000
33      0020000
34      0021000
35      0022000
36      0023000
37      0024000
38      0025000
39      0026000
40      0027000
41      0028000
42      0029000
43      0030000
44      0031000
45      0032000
46      0033000
47      0034000
48      0035000
49      0036000
50      0037000
51      0038000
52      0039000
53      0040000
54      0041000
55      0042000
56      0043000
57      0044000
58      0045000
59      0046000
60      0047000
61      0048000
62      0049000
63      0050000
64      0051000
65      0052000
66      0053000
67      0054000
68      0055000
69      0056000
70      0057000
71      0058000
72      0059000
73      0060000
74      0061000
75      0062000
76      0063000
77      0064000
78      0065000
79      0066000
80      0067000
81      0068000
82      0069000
83      0070000
84      0071000
85      0072000
86      0073000
87      0074000
88      0075000
89      0076000

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*** MEMBER U8P1

784	RCM(I,1)*RCRP(I,2)	0077000
794	RCM(I,2)*RCRP(I,3)	0078000
804	240 CONTINUE	0079000
814	241 CONTINUE	0080000
824	230 CONTINUE	0081000
834	IF(NB,EG,NBLK) GO TO 246	0082000
844	DO 245 I=1,NP2	0083000
854	RINV(I,I,NB)*RR(I+1,1)	0084000
864	245 CONTINUE	0085000
874	DO 247 I=1,NP2	0086000
884	DUHMY1(I,I)*RR(I+1,2)	0087000
894	247 CONTINUE	0088000
904	GO TO 249	0089000
914	249 CONTINUE	0090000
924	DO 244 I=2,NP1	0091000
934	DUHMY1(I,I-1)*AX(I)*RCRR(I-1,2)+AY(NP-1)*RCRR(I,1)*	0092000
944	1AH(I,NP-1)*RCRR(I,2)+CX(I)*RCRR(I+1,2)	0093000
954	248 CONTINUE	0094000
964	249 CONTINUE	0095000
974	215 CONTINUE	0096000
984	CALL MATINV(DUHMY1)	0097000
994	IF(NB,EG,NBLK) GO TO 275	0098000
1004	DO 260 J=1,NP2	0099000
1014	DO 260 I=1,NP2	0100000
1024	RINV(I,J,NB)=0.0	0101000
1034	DO 261 I=1,NP2	0102000
1044	RINV(I,J,NB)*RINV(I,J,NB)+DUHMY1(I,K)*RINV(K,J,NB)	0103000
1054	261 CONTINUE	0104000
1064	260 CONTINUE	0105000
1074	DO 270 J=1,NP2	0106000
1084	DO 270 I=1,NP2	0107000
1094	RINV(I,J,NB)=DUHMY1(I,J)	0108000
1104	270 CONTINUE	0109000
1114	275 CONTINUE	0110000
1124	295 CONTINUE	0111000
1134	RETURN	0112000
1144	END	0113000
		0114000

*** MEMBER MATINV

1*	SUMMATIVE MATINV(B)	
2*	PARAMETER MP21	0001000
3*	PARAMETER MP34=1	0002000
4*	REALDP R1(MP,MP)	0003000
5*	REALDP R2(MP,MP)	0004000
6*	MP1,MP21	0005000
7*	C9 110 I=1,MP1	0006000
8*	B1(I)=1.0/B(I,I)	0007000
9*	B(I,I)=1.0	0008000
10*	DO 112 J=1,MP	0009000
11*	B(I,J)=B(I,J)+B1(I)	0010000
12*	112 CONTINUE	0011000
13*	IP1=I+1	0012000
14*	DO 120 I=1,MP1,4	0013000
15*	B1(I)=B(I,I)	0014000
16*	120 CONTINUE	0015000
17*	DO 125 I=1,IP1,4	0016000
18*	B(I,I)=0.0	0017000
19*	125 CONTINUE	0018000
20*	DO 127 J=1,MP	0019000
21*	B2(J)=B(I,J)	0020000
22*	127 CONTINUE	0021000
23*	DO 135 I=1,IP1,4	0022000
24*	B(I,I)=0.0	0023000
25*	B(I,I)=B(I,I)-B1(I,I)+B2(J)	0024000
26*	135 CONTINUE	0025000
27*	140 CONTINUE	0026000
28*	B1(I)=B1(I)/4(MP,MP)	0027000
29*	B(MP,MP)=1.0	0028000
30*	DO 140 J=1,MP	0029000
31*	B(MP,J)=B(MP,J)+B1(I)	0030000
32*	140 CONTINUE	0031000
33*	DO 150 I=2,MP	0032000
34*	DO 155 I=2,I	0033000
35*	B1(I2)=B(I2,I)	0034000
36*	155 CONTINUE	0035000
37*	I=I+1	0036000
38*	DO 156 I=2,I,IM1	0037000
39*	B(I2,I)=J.0	0038000
40*	156 CONTINUE	0039000
41*	DO 157 J=1,MP	0040000
42*	B2(J)=B(I,J)	0041000
43*	157 CONTINUE	0042000
44*	I=I+1	0043000
45*	DO 160 I=2,I,IM1	0044000
46*	DO 160 J=1,MP	0045000
47*	B(I2,J)=B(I2,J)-B1(I2)+B2(J)	0046000
48*	160 CONTINUE	0047000
49*	150 CONTINUE	0048000
50*	RETURN	0049000
51*	END	0050000
		0051000

*** MEMBER BSM2

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1*      SUBROUTINE BSM2(X,FNR0P,A11,A1N,A21,A2N)          0001000
2*      PARAMETER N=21,N=21
3*      PARAMETER NBLK=N2,NBL=1+NBLK=1                      0002000
4*      PARAMETER NP=N+1,NPN=N+1                            0003000
5*      PARAMETER NP1=NP+1,NP2=NP+2,NP1NP=1,NP2NP=2        0004000
6*      REAL*8 RCRR,RINV,RINV1,RTILDA                      0005000
7*      COMMON/EVP/RINV(NP2,NP2,NBLK),RINV1(NP2,NP2,NBLK),RCRR(NP,3),    0006000
8*      1          RTILDA(NP2),F(NP,NP),NBSIZZ(NBLK),IS(NBLK),SLMF(NBLK),    0007000
9*      2          IE(NBLK),F11(NP),F1N(NP),F21(NP),F2N(NP),AX(NP),AY(NP),    0008000
10*      3          RR(NP,NP),CX(NP),CY(NP)                  0009000
11*      DIMENSION X(NP,NP)                                0010000
12*      DA 90 NB1,NBLK                                    0011000
13*      SUMF(NB)=0.0                                     0012000
14*      90 CONTINUE                                     0013000
15*      DA 95 NB1,NBLK                                    0014000
16*      DA 95 I=2,NP1                                    0015000
17*      SUMF(NB)=SUMF(NB)+ABS(F(I,IE(NB)-1))           0016000
18*      95 CONTINUE                                     0017000
19*      DA 96 NB1,NBLK                                    0018000
20*      IF(SUMF(NB).GT.1.0) GE 78 96                  0019000
21*      SUMF(NB)=1.0                                     0020000
22*      96 CONTINUE                                     0021000
23*      NSTARTS1                                      0022000
24*      DA 199 IT3=1,5                                 0023000
25*      DA 200 NB=NSTART,NBLK                         0024000
26*      ISP1=IS(NB)+1                                0025000
27*      IEM2=IE(NB)+2                                0026000
28*      DA 205 J=ISP1,IEM2                           0027000
29*      DA 205 I=2,NP1                                0028000
30*      X(I,J+1)=F(I,J)-AX(I)*X(I-1,J)-AY(J)*X(I,J-1)-BB(I,J)+ 0029000
31*      IX(I,J)=CX(I)*X(I+1,J))/CY(J)              0030000
32*      205 CONTINUE                                     0031000
33*      IF(NB.EQ.NBLK) GO TO 200                     0032000
34*      DA 522 IT1=1,10                               0033000
35*      J1=IE(NB)-1                                  0034000
36*      DA 215 I=2,NP1                                0035000
37*      RTILDA(I+1)=X(I,J1+1)-(F(I,J1)-AX(I)*X(I-1,J1)-AY(J1))+ 0036000
38*      1X(I,J1+1)=BB(I,J1)+X(I,J1)                -CX(I)*X(I+1,J1))/CY(J1) 0037000
39*      215 CONTINUE                                     0038000
40*      42=0.0                                       0039000
41*      DA 216 I=1,NP2                                0040000
42*      42=42*DARS(RTILDA(I))                      0041000
43*      216 CONTINUE                                     0042000
44*      A3=42/SLMF(NB)                                0043000
45*      IF(A3.LE.0.1) GE 19 230                      0044000
46*      DA 217 J=1,3                                 0045000
47*      DA 217 I=1,NP                                0046000
48*      RCRR(I,J)=0.0                                0047000
49*      217 CONTINUE                                     0048000
50*      DA 223 J=1,NP2                                0049000
51*      RCRR(J+1,2)=0.0..                            0050000
52*      DA 223 J=1,NP2                                0051000
53*      RCRR(J+1,2)=RCRR(J+1,2)+RTILDA(J)+RINV(J1,J,NB) 0052000
54*      223 CONTINUE                                     0053000
55*      IF(NB.EQ.1) GO TO 251                      0054000
56*      DA 225 J=2,NP1                                0055000
57*      RCOM(J,1)=0.0                                0056000
58*      DA 225 K=2,NP1                                0057000
59*      RCOM(J,1)=RCOM(J,1) + RCOM(K,2)*RINV(K=1,J=1,NB=1) 0058000
60*      225 CONTINUE                                     0059000
61*      DA 226 I=2,NP1                                0060000
62*      X(I,I,NB)=X(I,I,NB)+RCRR(I,I)              0061000
63*      226 CONTINUE                                     0062000
64*      251 CONTINUE                                     0063000
65*      CALL BSM3(X,IS(NB),IE(NB))                  0064000
66*      522 CONTINUE                                     0065000
67*      230 CONTINUE                                     0066000
68*      J1=IE(NB)-1                                  0067000
69*      DA 220 I=2,NP1                                0068000
70*      X(I,J1+1)=F(I,J1)-AX(I)*X(I-1,J1)-AY(J1)*X(I,J1-1)- 0069000
71*      BB(I,J1)*X(I,J1)-(X(I)+X(I+1,J1))/CY(J1) 0070000
72*      220 CONTINUE                                     0071000
73*      501 CONTINUE                                     0072000
74*      200 CONTINUE                                     0073000
75*      DA 300 NB1=N1,NBLK                          0074000
76*      NB=NBLK =NB1+1                                0075000
77*      ISP1=IS(NB)+1                                0076000

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*** MEMBER BSM2

```

78*      IEM2=IE(NR)=2          0078000
79*      J81EM2                  0079000
80*      IF(NB,EG,NBLK) GO TO 502 0080000
81*      DO 305 I=2,MP1          0081000
82*      X(I,J+1)=F(I,J)-AX(I)*X(I-1,J)-AY(J)*X(I,J+1)-BB(I,J)*
83*      IX(I,J)-CX(I)*X(I+1,J))/CY(J) 0082000
84*      305 CONTINUE            0083000
85*      502 CONTINUE            0084000
86*      DO 552 J=1,10           0085000
87*      IF(NB,EG,NBLK) GO TO 317 0086000
88*      J18IE(NB)=1             0087000
89*      DO 315 I=2,MP1          0088000
90*      RTILDA(I=1)*X(I,J+1)=(F(I,J1)-AX(I)*X(I-1,J1)-AY(J1)*
91*      IX(I,J1)-BB(I,J1)*X(I,J1)*CX(I)*X(I+1,J1))/CY(J1) 0089000
92*      315 CONTINUE            0090000
93*      GO TO 318              0091000
94*      317 CONTINUE            0092000
95*      DO 319 I=2,MP1          0093000
96*      RTILDA(I=1)*F(I,NP=1)=(AX(I)*X(I-1,NP=1)+AY(NP=1)*X(I,NP=2)*
97*      BB(I,NP=1)*X(I,NP=1)*CX(I)*X(I+1,NP=1)) 0094000
98*      319 CONTINUE            0095000
99*      318 CONTINUE            0096000
100*     A2=0,0                   0100000
101*     DO 316 J=1,MP2          0101000
102*     A2=A2*DABS(RTILLA(I    )) 0102000
103*     316 CONTINUE            0103000
104*     A3=A2/SUMF(NB)          0104000
105*     IF(A3.LE.ERROR) GO TO 300 0105000
106*     DO 320 J=1,3             0106000
107*     DO 320 J=1,MP             0107000
108*     RCAR(I,J)=0,0            0108000
109*     320 CONTINUE            0109000
110*     DO 324 J=1,MP2          0110000
111*     RCAR(J+1,2)=0,0          0111000
112*     DO 324 J=1,MP2          0112000
113*     RCAR(J+1,2)=RCBR(J+1,2)+RTILDA(J1)*RINV(J1,J,NB) 0113000
114*     324 CONTINUE            0114000
115*     IF(NB,EG,1) GO TO 551    0115000
116*     DO 325 J=2,MP1          0116000
117*     RCAR(J,1)=0,0            0117000
118*     DO 325 K=2,MP1          0118000
119*     RCAR(J,1)=RCGR(J,1)      +RCBR(K,2)*RINV1(K=1,J=1,NB=1) 0119000
120*     325 CONTINUE            0120000
121*     DO 326 I=2,MP1          0121000
122*     X(I,IS(NB))*X(I,IS(NB))+RCBR(I,1) 0122000
123*     326 CONTINUE            0123000
124*     551 CONTINUE            0124000
125*     CALL BSM3(X,IS(NB),IE(NB)) 0125000
126*     552 CONTINUE            0126000
127*     300 CONTINUE            0127000
128*     J18IE(I)                0128000
129*     DO 330 I=2,MP1          0129000
130*     RTILDA(I=1)*X(I,J+1)=(F(I,J1)-AX(I)*X(I-1,J1)-AY(J1)*
131*     IX(I,J1)-BB(I,J1)*X(I,J1)          -CX(I)*X(I+1,J1))/CY(J1) 0130000
132*     330 CONTINUE            0131000
133*     A2=0,0                   0132000
134*     DO 332 J=1,MP2          0133000
135*     A2=A2*DABS(RTILLA(I)) 0134000
136*     332 CONTINUE            0135000
137*     A3=A2/SUMF(1)          0136000
138*     IF(A3.LE.ERROR) GO TO 201 0137000
139*     NSTART=2                 0138000
140*     199 CONTINUE            0139000
141*     201 CONTINUE            0140000
142*     DO 350 J=2,MP1          0141000
143*     X(I,J)=(1,0=A21)*X(I,J)+A21*(X(2,J)=F21(J)) 0142000
144*     X(MP,J)=(1,0=A21)*X(MP,J)+A21*(X(MP-1,J)=F21(J)) 0143000
145*     350 CONTINUE            0144000
146*     DO 366 I=2,MP1          0145000
147*     X(I,I)=(1,0=A11)*X(I,I)+A11*(X(1,2)=F11(I)) 0146000
148*     X(I,NP)=(1,0=A11)*X(I,NP)+A11*(X(I,NP-1)=F11(I)) 0147000
149*     366 CONTINUE            0148000
150*     DO 371 J=2,MP1          0149000
151*     BB(2,J)=BB(2,J)+A1(2)*A21 0150000
152*     F(2,J)=F(2,J)+A1(2)*F21(J)+A21 0151000
153*     BB(MP-1,J)=BB(MP-1,J)+CX(MP-1)*A21 0152000

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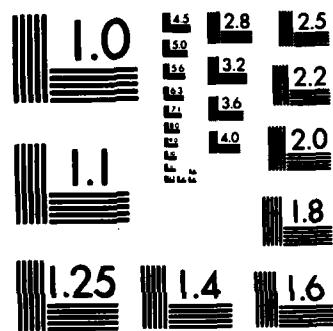
AD-A137 421 TROPICAL WEATHER SYSTEM AND OCEAN MODELING(U) SCIENCE 2/2
APPLICATIONS INC MCLEAN VA S CHANG 1983 SAI-83-1155
N00014-82-C-2306

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MICROCOPY RESOLUTION TEST CHART
NATIONAL BUREAU OF STANDARDS-1963-A

*** MEMBER H8P2

154*	F(NP+1,J)*F(NP+1,J)*CX(NP+1)*F2*(J)*A2H	0154000
155*	371 CONTINUE	0155000
156*	D2 372 IM2,NP1	0156000
157*	BB(I,2)*BB(I,2)*AV(2)*A11	0157000
158*	F(I,2)*F(I,2)*AV(2)*F11(I)*A11	0158000
159*	BB(I,NP+1)*BB(I,NP+1)*CY(NP+1)*A1N	0159000
160*	F(I,NP+1)*F(I,NP+1)*CY(NP+1)*F1N(I)*A1N	0160000
161*	372 CONTINUE	0161000
162*	RETURN	0162000
163*	END	0163000

*** MEMBER USM3

```

1+      SUBROUTINE USM3(X,ISS,IEE)          0001000
2+      PARAMETER NBLK2,NBLK1           0002000
3+      PARAMETER NBLK2,NBLK1+NBLK1   0003000
4+      PARAMETER NP1NP01,NP1NP02   0004000
5+      PARAMETER NP2NP01,NP2NP02,NP1NP01,NP2NP02   0005000
6+      REAL A0,RGR,RILV,LINV1,RTILDA   0006000
7+      DIMENSION X(NP,NP)           0007000
8+      COMMON/EVR/RINV(NP2,NP2,NBLK),PINV1(NP2,NP2,NBLK1),RCOR(NP,3),
9+                  RTILDA(NP2),P(NP,NP),NBSIZZ(NBLK),IS(NBLK),SLMP(NBLK), 0008000
10+                 IE(NBLK),F1(NP),F11(NP),F21(NP),F2M(NP),AX(-P),AY(NP), 0009000
11+                 90(NP,NP),CX(NP),CY(NP)   0010000
12+      DO 135 I=1,NP1           0011000
13+      X(I,ISS+1)=X(I,ISS+1)+RCOR(I,2)   0012000
14+ 135 CONTINUE           0013000
15+      ISS=ISS+1           0014000
16+      IEN2=IEE+2           0015000
17+      DO 140 J=1,NP1           0016000
18+      DO 145 I=2,NP1           0017000
19+      RCOR(I,3)=(-AX(J))+RCOR(I-1,2)-AY(J)+RCOR(I,1)-90(I,J)+ 0018000
20+      1NCOR(I,2)=CX(I)+RCOR(I+1,2))/CY(J)   0019000
21+ 145 CONTINUE           0020000
22+      DO 150 I=2,NP1           0021000
23+      X(I,J+1)=X(I,J+1)+RCOR(I,3)   0022000
24+      RCOR(I,1)=RCOR(I,2)   0023000
25+      RCOR(I,2)=RCOR(I,3)   0024000
26+ 150 CONTINUE           0025000
27+ 160 CONTINUE           0026000
28+      RETURN           0027000
29+      END           0028000
30+                          0029000

```

<<SPLIT OCEAN1,SOURCE1,PRINT,SEG

```

10      PROGRAM OCEAN          0001000
20      PARAMETER NO21,NO22  0002000
30      PARAMETER N12NO1,N22NO2 0003000
40      PARAMETER NO22NO1+N12NO1 0004000
50      DIMENSION DATA1(NC),DATA2(NO),DATA3(NO) 0005000
60      COMMON//BLK/VRI(1,N1),VT1(N,N1),VI(N1,N1),VR2(N,N1), 0006000
70      VT2(N,N1),B2(N1,N1),VR3(N,N1),VT3(N,N1), 0007000
80      B3(N1,N1),P(N1,N1),VZ(N1,N1) 0008000
90      EQUIVALENCE (DATA1,VR1),(DATA2,VR2),(DATA3,VR3) 0009000
100     DATA DATA1/NO21,,/DATA2/NO22,,/DATA3/NO23/ 0010000
110     COMMON//TH0/RH0(NM),B2(N1),DR1(N1),DR2(N),Z1(N),Z2(N1),CZ1(N1),DZ2(N) 0011000
120     COMMON//TH1/RH1(NM),B1(N1),DR1(N1),DR2(N),Z1(N),Z2(N1),CZ1(N1),DZ2(N) 0012000
130     COMMON//FOR/DELT,XTIME,ITIME,ISTEP,ISNO,ITAPE,TBV 0013000
140     CALL INRDUMP 0014000
150 100 FORMAT(1A) 0015000
160     READ(5,100)ITIME 0016000
170     READ(5,100)ITER 0017000
180     READ(5,100)IOUT 0018000
190     READ(5,100)ISNO 0019000
200     ITER=0 0020000
210     READ(5,100)ITAPE 0021000
220     CALL INIT 0022000
230     IF(ITIME,EG,0)GO TO 10 0023000
240 C 0024000
250 C CONTINUED INTEGRATION FROM A HISTORY TAPE 0025000
260 C 0026000
270     READ(1)ITIME,DATA1,DATA2,P 0027000
280     GO TO 20 0028000
290 10 CALL START 0029000
300 20 ITIME=ITIME+3000. 0030000
310 C 0031000
320 C PRINT BLT INITIAL FIELDS 0032000
330 C 0033000
340     CALL PUTBT 0034000
350     IF(ITER,EG,0)STOP 0035000
360     GO TO ISTEP1,ITER 0036000
370 C 0037000
380 C COMPUTE HYDROSTATIC PRESSURE AND DIAGNOSE VERTICAL VG 0038000
390 C 0039000
400 CALL UP 0040000
410 C 0041000
420 C COMPUTE ALL INVISCID TERMS 0042000
430 C 0043000
440 CALL ADVECT 0044000
450 C 0045000
460 C COMPUTE VISCOUS TERMS 0046000
470 C 0047000
480 CALL DIFF 0048000
490 C 0049000
500 C MARCHING IN TIME 0050000
510 C FIRST TIME STEP IS FORWARD IF START IS CALLED 0051000
520 C 0052000
530 ISTEP,EG,1,AND,ITIME,EG,0)DELTA0.9=DELT 0053000
540 CALL FWDAD 0054000
550 ISTEP,EG,1,AND,ITIME,EG,0)DELTA0.9=DELT 0055000
560 C 0056000
570 C DEFINE BOUNDARY VALUES FOR VELOCITY 0057000
580 C 0058000
590 CALL BOUNDV 0059000
600 C 0060000
610 C CHECK IF DELT IS STABLE 0061000
620 CC 0062000
630 C 0063000
640 CALL CHECK 0064000
650 XTIME=ITIME+DELT 0065000
660 ITIME=ITIME/3000. 0066000
670 C 0067000
680 C PRINT BLT RESULTS EVERY IOUT STEPS 0068000
690 C 0069000
700 C IF(MOD(ISTEP,10LT),EG,0)CALL PUTBT 0070000
710 C 0071000
720 C WRITE HISTORY TAPE EVERY ITAPE STEPS 0072000
730 C IF(MOD(ISTEP,ITAPE),EG,0)WRITE(2)ITIME,DATA1,DATA2,P 0073000
740 GO CONTINUE 0074000
750 STOP 0075000
760 END 0076000

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SEE MEMBER INIT

```

10      SUBROUTINE INIT
20      PARAMETER NB21,NB21
30      PARAMETER N10N=1,N20N=2,N10N=1,N20N=2
40      COMMON/DNE/VR1(N,N1),VT1(N,N1),R1(N1,N1),VR2(N,N1),
50      1 VT2(N,N1),R2(N1,N1),VR3(N,N1),VT3(N,N1),
60      2 R3(N1,N1),P(N1,N1),VZ(N1,N)
70      COMMON/TNO/R1(N),R2(N1),DR1(N1),DR2(N),Z1(N),Z2(N1),DZ1(N1),DZ2(N)
80      COMMON/TNR/RND,BNDR(N1),BV2(N),ALPHA,BND,BNDR,CORI,E,NK(N),ZK(N)
90      COMMON/FOR/DELT,XTIME,ITIME,ISTEP,IBNG,ITAPE,TBV
100 C
110 C      INITIALIZE ALL DEPENDENT VARIABLES AND CONSTANTS
120 C
130 C
140 C      ALPHA IS THE NONDIMENSIONAL SMOOTHING COEF.
150 C      FOR TIME SMOOTHING IN SUBROUTINE FNRAD
160 C
170 C      DELT=999,
180 C      ALPHA=0.30
190 C      BN999,
200 C      LAT030
210 C      CORI=2.072732E-5+8IN(LAT+3.14159/180.)
220 C
230 C      DEFINE RADII AT GRID POINTS AND ALL GRID INTERVALS
240 C
250      DO 10 I=1,N1
260      10 DR1(I)=26.E3
270      R1(I)=0.
280      DO 20 I=2,N
290      20 R1(I)=R1(I-1)+DR1(I-1)
300      DO 30 J=1,N1
310      30 R2(J)=0.5*(R1(J)+R1(J+1))
320      DR2(1)=2.0*(R2(1)-R1(1))
330      DR2(N)=2.0*(R1(N)-R2(N1))
340      DO 40 I=2,N1
350      40 DR2(I)=R2(I)-R2(I-1)
360      CMAX=MAX(DR1)
370      DRMAX=DR1(CMAX)
380 C
390 C      DEFINE ALL DZ'S
400 C
410      DO 100 J=1,N
420      100 Z1(J)=Z1(J-1)+200.E3
430      DO 110 J=1,N1
440      110 DZ1(J)=Z1(J+1)-Z1(J)
450      DZ2(1)=2.0*(Z1(1)-Z1(1))
460      Z2(1)=0.5*DZ2(1)
470      DO 120 J=2,N1
480      120 DZ2(J)=Z2(J-1)-DZ2(J)
490      120 Z2(J)=Z2(J-1)+DZ2(J)
500      120 DZ2(N)=Z2(Z(N))-Z2(N-1)=0.5*DZ2(N1))
510 C
520 C      A AND B ARE CONSTANT USED IN SUBROUTINE BOUNDV
530 C      FOR CONSTANT DIV. AND VORT. CONDITIONS
540 C
550      BN048R1(N1)/R1(N)
560      BN008R1(N1)+R1(N))+DR1(N1)/(C(R1(N1)+R1(N2))+R1(N)+DR1(N2))
570      BN008R1.
580 C
590 C      DEFINE CONDUCTIVITY RELATED CONSTANTS
600 C
610      RN001,
620      DO 130 J=1,N
630      130 BV2(J)=1.E-6
640      TBV00,
650      DO 135 J=1,N
660      135 TBV=MAX(TBV,BV2(J))
670      TBV1=.80RT(TBV)
680 C
690 C      DEFINE HORIZONTAL AND VERTICAL DIFFUSION COEFFICIENTS
700 C
710      C00FHN00.002*DR1(1)*=2/DELT
720      C00FZ00.001*DZ1(1)*=2/DELT
730      DO 140 I=1,N1
740      140 NK(I)=C00FHN00.005.*EXP(-FLOAT(N1-I)/7.0)
750      DO 150 J=1,N1
760      150 ZK(J)=C00FZ00.005.*EXP(-FLOAT(J-1)/5.0)*EXP(-FLOAT(N1-J)/5.0))
770      RETURN
780      END

```

*** MEMBER START

```

10      SUBROUTINE START
20      PARAMETER M021,L021
30      PARAMETER M10N=1,M20N=2,M10N=1,M20N=2
40      COMMON/CNE/VRI(N,N),VT1(N,N),V1(N,N),VR2(N,N),
50      VT2(N,N),V2(N,N),VR3(N,N),VT3(N,N),
60      V3(N,N),P(N,N),V2(N,N)
70      COMMON/TWA/R1(N),R2(N),R3(N),DR1(N),DR2(N),Z1(N),Z2(N),GZ1(N),DZ2(N)
80      COMMON/TWA/RHO,ELGR(N),BV2(N),ALPHA,BNDA,BNDB,CORI,G,MK(N),ZK(N)
90      PARAMETER M020P=N1+M1+N1
100     DIMENSION DATA1(NC),DATA2(NC)
110     EQUIVALENCE (DATA1,VR1),(DATA2,VR2)

120 C
130 C
140 C
150 I1I1
160 I1I2I3
170 I2I0I2+1
180 I20001
190 DR16N0,0002
200 ON 10 I0I1I1,I1I2
210 10 S1(I,N1)=BMA0+C08*(PL0AT(I=I11)/0.+3.14159)*G/RHO
220 00 30 I0I21,I22
230 30 S1(I,N1)=S1(I12,N1)*EXP(-PL0AT(I=I21+1)/0.)
240 02 40 J01,N2
250 FACT=EXP(PL0AT(J=41)/0.)
260 00 40 I01,N1
270 40 S1(I,J)=S1(I,N1)*FACT

280 C
290 C
300 C
310 C
320 C
330 C
340 C
350 C
360 C
370 C
380 C
390 C
400 C
410 C
420 C
430 C
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*** MEMBER UP

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10      SUBROUTINE UP          0001000
20      PARAMETER N021,L021  0002000
30      PARAMETER N10N01,N20N02,N10N01,N20N02  0003000
40      COMMON/SNE/VR1(P,1),VT1(N,N1),B1(N1,N1),VR2(N,N1),  0004000
50      1           VT2(P,N1),B2(N1,N1),VR3(N,N1),VT3(N,N1),  0005000
60      2           B3(N1,N1),P(N1,N1),VZ(N1,N1)  0006000
70      COMMON/TNS/R1(N),S2(N1),DR1(N1),DR2(N),Z1(N),Z2(N1),DZ1(N1),DZ2(N)  0007000
80      COMMON/THR/RNG,SNEA(N1),BV2(N),ALPHA,BNDA,BNDB,CBR1,S,NK(N),ZK(N)  0008000
90      PARAMETER N10N01,N20N02  0009000
100     DATA VZ/N10N02/  0010000
110 C
120 C             PRESSURE IS OBTAINED HYDROSTATICALLY FROM 0  0011000
130 C
140     DD 10 JN1,N1  0012000
150     10 P(I,1)=0,5*RN0*DZ2(I)+B2(I,1)  0013000
160     DC 20 JN2,N1  0014000
170     DD 20 JN1,N1  0015000
180     20 P(I,J)=P(I,J-1)+0,5*RN0*DZ2(J)+(B2(I,J)+B2(I,J-1))  0016000
190 C
200 C             DIAGNOSE VERTICAL VELOCITY BY CONTINUITY EQUATION  0017000
210 C
220     DD 30 JN2,N1  0018000
230     DD 30 JN1,N1  0019000
240     30 VZ(I,J)=VZ(I,J-1)+DZ1(J)*(R1(I+1)+VR2(I+1,J-1)-R1(I)*VR2(I,J-1))  0020000
250     1           /(DR1(I)+0,5*(R1(I+1)+R1(I)))  0021000
260     RETURN  0022000
270     END  0023000

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ONE MEMBER BOUNDARY

```

10      SUBROUTINE BOUNDARY          0001000
20      PARAMETERS N021,N021          0002000
30      PARAMETERS N10N01,N20N02,N10N01,N20N02          0003000
40      COMMON/BNE/VR1(N,N1),VT1(N,N1),B1(N1,N1),VR2(N,N1),          0004000
50      VT2(N,N1),B2(N1,N1),VR3(N,N1),VT3(N,N1),          0005000
60      1      B3(N1,N1),P(N1,N1),V2(N1,N1)          0006000
70      2      COMMON/TAB/R1(N),R2(N1),DR1(N1),DR2(N),Z1(N),Z2(N1),DZ1(N1),DZ2(N) 0007000
80      COMMON/TAB/RND,BNDL(N1),BV2(N),ALPHA,BNDL,BND0,BND1,G,MK(N),ZK(N) 0008000
90 C
100 C          LATERAL BOUNDARY FOR TANGENTIAL AND RADIAL VELOCITIES 0010000
110 C          ASSUMING CONTINUOUS VORTICITY AND DIVERGENCE          0011000
120 C
130      DD 10 J\$1,N1          0013000
140      VR2(N,J)=BNDL=VR2(N1,J)+BND0=(R1(N1)+VR2(N1,J)-R1(N2)+VR2(N2,J)) 0014000
150      10 VT2(N,J)=BNDL=VT2(N1,J)+BND0=(R1(N1)+VT2(N1,J)-R1(N2)+VT2(N2,J)) 0015000
160      RETURN          0016000
170      END          0017000

```

*** MEMBER DIFF

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10      SUBROUTINE DIFF          0001000
20
30      COMPLETE THE DIFFUSION TERMS 0002000
40
50      PARAMETER N021,N021          0003000
60      PARAMETER N10M=1,P20M=2,N10N=1,N20N=2          0004000
70      COMMON/ONE/VR1(N,N1),VT1(N,N1),B1(N,N1),VR2(N,N1),          0005000
80      VT2(N,N1),B2(N,N1),VR3(N,N1),VT3(N,N1),          0006000
90      B3(N,N1),P(N,N1),VZ(N,N1)          0007000
100     COMMON/TMR/RA1(N),R2(N),DR1(N),DR2(N),Z1(N),Z2(N),DZ1(N),DZ2(N) 0008000
110     COMMON/TMR/RM0,RM0(N),BV2(N),ALPHA,BNDA,BNDB,CORI,G,MK(V),ZK(N) 0009000
120     DIMENSION VR(N,N1),VT(N,N1),B(N,N1)          0010000
130     EQUIVALENCE (VR,VR1),(VT,VT1),(B,B1)          0011000
140
150      HORIZONTAL DIFFUSION OF RADIAL VELOCITY 0012000
160
170      DD 10 JN1,N1          0013000
180      DD 10 JN2,N1          0014000
190      10 VR3(I,J)=VR3(I,J)+MK(I)*((VR(I+1,J)-VR(I,J))/DR1(I)) 0015000
200      1      -(VR(I,J)-VR(I-1,J))/DR1(I-1))/DR2(I)-VR(I,J)/(R1(I)*R1(I)) 0016000
210      2      +0.5*((VR(I+1,J)-VR(I,J))/(DR1(I)*R2(I)))          0017000
220      3      +(VR(I,J)-VR(I-1,J))/(DR1(I-1)*R2(I-1)))          0018000
230
240      HORIZONTAL DIFFUSION OF TANGENTIAL VELOCITY 0019000
250
260      DD 20 JN1,N1          0020000
270      DD 20 JN2,N1          0021000
280      20 VT3(I,J)=VT3(I,J)+MK(I)*((VT(I+1,J)-VT(I,J))/DR1(I)) 0022000
290      1      -(VT(I,J)-VT(I-1,J))/DR1(I-1))/DR2(I)-VT(I,J)/(R1(I)*R1(I)) 0023000
300      2      +0.5*((VT(I+1,J)-VT(I,J))/(DR1(I)*R2(I)))          0024000
310      3      +(VT(I,J)-VT(I-1,J))/(DR1(I-1)*R2(I-1)))          0025000
320
330      HORIZONTAL DIFFUSION OF B 0026000
340
350      DD 60 JN1,N1          0027000
360      DD 60 JN2,N2          0028000
370      60 B3(I,J)=B3(I,J)+MK(I)*(((B(I+1,J)-B(I,J))/DR2(I+1) 0029000
380      1      -(B(I,J)-B(I-1,J))/DR2(I))/DR1(I))          0030000
390      2      +0.5*((B(I+1,J)-B(I,J))/(DR2(I+1)*R1(I+1)))          0031000
400      3      +(B(I,J)-B(I-1,J))/(DR2(I)*R1(I)))          0032000
410      DD 70 JN1,N1          0033000
420      70 B3(I,J)=B3(I,J)+MK(I)*((B(2,J)-B(1,J))/(DR2(2)*DR1(1)) 0034000
430      1      +0.5*(B(2,J)-B(1,J))/(DR2(2)*R1(2)))          0035000
440      DD 80 JN1,N1          0036000
450      80 B3(N1,J)=B3(N1,J)+MK(N1)*((-B(N1,J)+B(N2,J))/(DR2(N1)*DR1(N1)) 0037000
460      1      -(B(N1,J)-B(N2,J))/(DR2(N1)*R1(N1)))          0038000
470
480      VERTICAL DIFFUSION OF RADIAL VELOCITY 0039000
490
500      DD 90 JN2,N2          0040000
510      DD 90 JN2,N1          0041000
520      90 VR3(I,J)=VR3(I,J)+ZK(J)*((VR(I,J+1)-VR(I,J))/DZ2(J+1) 0042000
530      1      -(VR(I,J)-VR(I,J-1))/DZ2(J))/DZ1(J)          0043000
540      DD 100 JN2,N1          0044000
550      100 VR3(I,J)=VR3(I,J)+ZK(1)*((VR(I,2)-VR(I,1))/(DZ2(2)*DZ1(1)) 0045000
560      DD 110 JN2,N1          0046000
570      110 VR3(I,N1)=VR3(I,N1)+ZK(N1)*(-VR(I,N1)+VR(I,N2))/(DZ2(N1)*DZ1(N1)) 0047000
580
590      VERTICAL DIFFUSION OF TANGENTIAL VELOCITY 0048000
600
610      DD 120 JN2,N2          0049000
620      DD 120 JN2,N1          0050000
630      120 VT3(I,J)=VT3(I,J)+ZK(J)*((VT(I,J+1)-VT(I,J))/DZ2(J+1) 0051000
640      1      -(VT(I,J)-VT(I,J-1))/DZ2(J))/DZ1(J)          0052000
650      DD 130 JN2,N1          0053000
660      130 VT3(I,J)=VT3(I,J)+ZK(1)*((VT(I,2)-VT(I,1))/(DZ2(2)*DZ1(1)) 0054000
670      DD 140 JN2,N1          0055000
680      140 VT3(I,N1)=VT3(I,N1)+ZK(N1)*(-VT(I,N1)+VT(I,N2))/(DZ2(N1)*DZ1(N1)) 0056000
690

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ROW NUMBER DIFF

70+ C		0070000
71+ C	VERTICAL DIFFUSION OF B	0071000
72+ C		0072000
73+ 0F 160 J#2,N2		0073000
74+ 09 160 I#1,"1		0074000
75+ 160 B3(I,J)=B3(I,J)+ZK(J)*((B1(I,J+1)-B(I,J))/DZ2(J+1))		0075000
76+ 1 = (B(I,J)-B(I,J-1))/DZ2(J)/DZ1(J)		0076000
77+ 09 170 I#1,"1		0077000
78+ 170 B3(I,1)=B3(I,1)+ZK(1)*(B(I,2)-B(I,1))/(DZ2(2)*DZ1(1))		0078000
79+ 09 180 I#1,"1		0079000
80+ 180 B3(I,N1)=B3(I,N1)+ZK(N1)*(=B(I,N1)+B(I,N2))/(DZ2(N1)*DZ1(N1))		0080000
81+ RETURN		0081000
82+ END		0082000

*** MEMBER FWRD

```

10      SUBROUTINE FWRD
20      PARAMETER M1=21,M2=21
30      PARAMETER M1=M1+P2=2,N1=N1+1,N2=N2+2
40      PARAMETER ND=2*P+1+N1+N2
50      COMMON/ONE/DATA1(ND),DATA2(ND),DATA3(ND),P(M1,N1),V2(M1,N)
60      COMMON/TMR/RHO,RH0R(N1),BV2(N),ALPHA,BNDA,BNDB,C0R1,G,MK(N),ZK(N)
70      COMMON/PBR/DEL,T,XTIME,ITIME,ISTEP,ISM0,ITAPE,TBV
80
90 C          REPLACE DATA3 WITH THE NEW VALUES
100 C
110      DO 10 I=1,ND
120      10 DATA3(I)=DATA1(I)+2.*GELT*DATA3(I)
130 C          TIME SMOOTHING
140 C
150      IF(MOD(ISTEP,ISM0).NE.0)GO TO 30
160      DO 20 I=1,ND
170      20 DATA2(I)=DATA2(I)+(DATA1(I)+DATA3(I)-2.*DATA2(I))/ALPHA
180      30 CONTINUE
190
200 C          FORWARD MARCHING
210 C
220      DO 40 I=1,ND
230      40 DATA1(I)=DATA2(I)
240      DO 50 I=1,ND
250      50 DATA2(I)=DATA3(I)
260
270 C          ZERO OUT DATA3 FOR NEXT STEP
280 C
290      DO 60 I=1,ND
300      60 DATA3(I)=0.
310      RETURN
320
330      END

```

SEE MEMBER CHECK

```

1*      SUBROUTINE CHECK
2*      PARAMETER M=21,N=21
3*      PARAMETER M1=M-1,N1=N-1,M2=M-2,N2=N-2
4*      COMMON/DNE/VR1(M,N1),VT1(M,N1),B1(M1,N1),VR2(M,N1),
5*      1          VT2(M,N1),B2(M1,N1),VR3(M,N1),VT3(M,N1),
6*      2          B3(M1,N1),P(M1,N1),VZ(M1,N)
7*      COMMON/TB9/R1(M),F2(M1),DR1(M1),DR2(M),Z1(N),Z2(N1),CZ1(N1),DZ2(N)
8*      COMMON/FOR/DELT,XTIME,ITIME,ISTEP,ISRC,ITAPE,TBY
9*      DIMENSION WORK1(M),WORK2(N)
10*     DD 10 J=1,N1
11*     DD 20 I=1,M
12*     20 WORK1(I)=WORK2(I)/AMAX1(1.,VR2(I,J))
13*     MINBMINMAC(WORK1)+1
14*     DT=AMIN1(MIN)+0.9
15*     DT=AMIN1(DT,DELT)
16*     10 CONTINUE
17*     DT=AMIN1(DT,TBY)
18*     IF(DT.GE.DELT)RETURN
19*     DELT=0.75*DELT
20*     PRINT 100,DELT
21*   100 FORMAT(////////,'*****DELT IS CHANGED TO',1PE11.2,'*****')
22*     RETURN
23*     END

```

*** MEMBER ADVECT

```

16      SUBROUTINE ADVECT          0001000
20  C                                     0002000
30  C      COMPLETE THE ADVECTIVE TERMS 0003000
40  C                                     0004000
50  C                                     0005000
60  C      PARAMETER M821,M821          0006000
70  C      COMMON/CNE/VR1(M,N1),VT1(M,N1),B1(M,N1),VR2(M,N1), 0007000
80  C      VT2(M,N1),B2(M,N1),VR3(M,N1),VT3(M,N1), 0008000
90  C      B3(M,N1),B(M,N1),VZ(M,N1) 0009000
100 C      COMMON/TNO/R1(M),R2(M),DR1(M),DR2(M),Z1(N),Z2(N1),DZ1(N1),DZ2(N) 0010000
110 C      COMMON/THR/RW0,RHGR(N1),BV2(N),ALPH0,BND4,BNDB,CORI,S,HK(M),ZK(N) 0011000
120 C      DIMENSION VR(M,N1),VT(M,N1),B(M,N1) 0012000
130 C      EQUIVALENCE (VR,VR2),(VT,VT2),(B,B2) 0013000
140 C                                     0014000
150 C                                     0015000
160 C                                     0016000
170 DO 10 J=1,N1 0017000
180 DO 10 I=2,M1 0018000
190 10 VR3(I,J)=0.25*((VR(I,J)+VR(I-1,J))+(VR(I,J)-VR(I-1,J))/DR1(I-1) 0019000
200 1       +(VR(I+1,J)+VR(I,J))+(VR(I+1,J)-VR(I,J))/DR1(I)) 0020000
210 2       +VR3(I,J) 0021000
220 C                                     0022000
230 C      HORIZONTAL ADVECTION FOR TANGENTIAL VELOCITY 0023000
240 C                                     0024000
250 DO 20 J=1,N1 0025000
260 DO 20 I=2,M1 0026000
270 20 VT3(I,J)=0.25*((VR(I,J)+VR(I-1,J))+(VT(I,J)-VT(I-1,J))/DR1(I-1) 0027000
280 1       +(VR(I+1,J)+VR(I,J))+(VT(I+1,J)-VT(I,J))/DR1(I)) 0028000
290 2       +VT3(I,J) 0029000
300 C                                     0030000
310 C      HORIZONTAL ADVECTION FOR BUBSTANCY 0031000
320 DO 30 J=1,N1 0032000
330 DO 30 I=2,M2 0033000
340 60 B3(I,J)=B3(I,J)+0.5*(VR(I,J)+B(I,J)+B(I-1,J))/DR2(I) 0034000
350 1       +VR(I+1,J)+B(I+1,J)+B(I,J))/DR2(I+1) 0035000
360 DO 70 J=1,N1 0036000
370 70 B3(I,J)=B3(I,J)+0.5*VR(2,J)*(B(2,J)+B(1,J))/DR2(2) 0037000
380 DO 80 J=1,N1 0038000
390 80 B3(M1,J)=B3(M1,J)+0.5*VR(M1,J)*(B(M1,J)+B(M2,J))/DR2(M1) 0039000
400 C                                     0040000
410 C      VERTICAL ADVECTION FOR RADIAL VELOCITY 0041000
420 C                                     0042000
430 DO 90 J=2,N2 0043000
440 DO 90 I=2,M1 0044000
450 90 VR3(I,J)=VR3(I,J)+0.25*((VZ(I-1,J)+VZ(I,J))+(VR(I,J)-VR(I,J-1)) 0045000
460 1       /DZ2(J)+(VZ(I,J+1)+VZ(I-1,J+1))+(VR(I,J+1)-VR(I,J))) 0046000
470 2       /DZ2(J+1)) 0047000
480 DO 95 I=2,M1 0048000
490 95 VR3(I,I)=VR3(I,I)+0.25*(VZ(I,2)+VZ(I-1,2))+(VR(I,2)-VR(I,1)) 0049000
500 1       /DZ2(2) 0050000
510 DO 96 I=2,M1 0051000
520 96 VR3(I,N1)=VR3(I,N1)+0.25*(VZ(I-1,N1)+VZ(I,N1))+(VR(I,N1)-VR(I,N2)) 0052000
530 1       /DZ2(N1) 0053000
540 C                                     0054000
550 C      VERTICAL ADVECTION FOR TANGENTIAL VELOCITY 0055000
560 C                                     0056000
570 DO 100 J=2,M2 0057000
580 DO 100 I=2,M1 0058000
590 100 VT3(I,J)=VT3(I,J)+0.25*((VZ(I-1,J)+VZ(I,J))+(VT(I,J)-VT(I,J-1)) 0059000
600 1       /DZ2(J)+(VZ(I,J+1)+VZ(I-1,J+1))+(VT(I,J+1)-VT(I,J))) 0060000
610 2       /DZ2(J+1)) 0061000
620 DO 105 I=2,M1 0062000
630 105 VT3(I,I)=VT3(I,I)+0.25*(VZ(I,2)+VZ(I-1,2))+(VT(I,2)-VT(I,1)) 0063000
640 1       /DZ2(2) 0064000
650 DO 106 I=2,M1 0065000
660 106 VT3(I,N1)=VT3(I,N1)+0.25*(VZ(I,N1)+VZ(I-1,N1))+(VT(I,N1)-VT(I,N2)) 0066000
670 1       /DZ2(N1) 0067000
680 C                                     0068000
690 C      VERTICAL ADVECTION FOR B 0069000
700 C                                     0070000
710 DO 150 J=2,M2 0071000
720 DO 150 I=1,M1 0072000
730 150 B3(I,J)=B3(I,J)+0.5*(VZ(I,J)+B(I,J)+B(I,J-1))/DZ2(J) 0073000
740 1       +VZ(I,J+1)+B(I,J+1)+B(I,J))/DZ2(J+1) 0074000
750 DO 160 I=1,M1 0075000
760 160 B3(I,1)=B3(I,1)+0.5*VZ(I,2)*(B(2,I,2)+B(1,I,1))/DZ2(2) 0076000
770 1       DO 170 I=1,M1 0077000

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*** MEMBER ADVECT

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78a    170 B3(I,N1)=B3(I,N1)+0.5*(VZ(I,N1)*(B(I,N1)-B(I,N2))/C22(N1))      0078000
79a  C
80a  C           INERTIA TERMS FOR HORIZONTAL MOMENTUM
81a  C
82a      DR 110 J=1,N1
83a      DR 110 I=2,M1
84a      VR3(I,J)*VR3(I,J)*VT(I,J)*(VT(I,J)/P1(I)+C8RI)
85a      110 VT3(I,J)*VT3(I,J)*VR(I,J)*(VT(I,J)/P1(I)+C8RI)
86a  C
87a  C           PRESSURE GRADIENT FORCE
88a  C
89a      DD 120 J=1,N1
90a      DD 120 I=2,M1
91a      120 VR3(I,J)*VP3(I,J)*(P(I,J)-P(I+1,J))/(RH0*DR2(I))
92a  C
93a  C           STRATIFICATION TERM
94a  C
95a      DD 130 J=1,N1
96a      DD 130 I=1,M1
97a      130 B3(I,J)=B3(I,J)+0.5*(VZ(I,J)*BV2(J)*VZ(I,J+1)*BV2(J+1))
98a      RETURN
99a      END

```

*** MEMBER PUTOUT

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1+      SUBROUTINE PUTOUT
2+      PARAMETER M021,M021
3+      PARAMETER M10M=1,P2BM=2,N10M=1,N2BM=2
4+      COMMON/BNE/VR1(P,N1),VT1(M,N1),R1(M,L1),VR2(M,N1),
5+      1          VT2(M,N1),R2(M,N1),VR3(M,N1),VT3(M,N1),
6+      2          R3(M,N1),P(M,N1),VZ(M,L1)
7+      COMMON/TMA/R1(M),R2(M),DR1(M1),DR2(M),Z1(M),Z2(M1),DZ1(M1),DZ2(M)
8+      COMMON/TMR/RMC,RHCR(N1),BV2(N),ALPHA,BNDA,BNDB,CORI,G,HR(M),ZK(N)
9+      COMMON/FAR/DELT,XTIME,ITIME,ISTEP,ISMO,ITAPE,TBV
10+ C
11+ C      THIS SUBROUTINE PRINT OUT FIELDS FOR A GLICK LOOK
12+ C
13+      DIMENSION IDUM(M,N)
14+      700 FORMAT(//,,,' HORIZONTAL VELOCITY (CM/S) AT T0'',16.1 M')
15+      705 FORMAT(//,,,' TANGENTIAL VELOCITY (CM/S) AT T0'',16.1 M')
16+      710 FORMAT(//,,,' VERTICAL VELOCITY (CM/S) AT T0'',16.1 M')
17+      715 FORMAT(//,,,' DIVERGENCE FIELD (+,-001) AT T0'',16.1 M')
18+      725 FORMAT(//,,,' PRESSURE (+10 DYN/CM**2) AT TIME0'',16.1 M') 0RI,
19+      720 FORMAT(1H1,/,,,,' OUTPUT AT TIME 0'',16.1 M' 0RI,
20+      1 F8.2,' DAY ISTEP 0'',17,' ++++++-----+')
21+      DAY=XTIME/86400.+0.0001
22+      PRINT 720,ITIME,DAY,ISTEP
23+      DD 10 J81,N1
24+      DD 10 I81,M
25+      10 IDUM(I,J)BVR2(I,J)
26+      PRINT 700,ITIME
27+      CALL "AP(IDUM,R1,Z2,M1)
28+      DR 20 J81,N1
29+      DR 20 I81,M
30+      20 IDUM(I,J)BVT2(I,J)
31+      PRINT 705,ITIME
32+      CALL "AP(IDUM,R1,Z2,M1)
33+      DD 30 J81,N
34+      DD 30 I81,M1
35+      30 IDUM(I,J)BVZ(I,J)
36+      PRINT 710,ITIME
37+      CALL "AP(IDUM,R2,Z1,M1)
38+      DR 30 J81,N1
39+      DD 30 I81,M1
40+      40 IDUM(I,J)=R2(I,J)+1.E3
41+      PRINT 715,ITIME
42+      CALL "AP(IDUM,R2,Z2,M1,M1)
43+      DR 30 J81,N1
44+      DD 30 I81,M1
45+      50 IDUM(I,J)=P(I,J)+1.E-1
46+      PRINT 725,ITIME
47+      CALL "AP(IDUM,R2,Z2,M1)
48+      RETURN
49+      END

```

ONE PICTURE MAP

1.	ROUTINE MAP(A,F,Z,MM,NN)	0001000
2.	PARAMETER MM21,LB21	0002000
3.	CIMENSIJA R(MM),Z(NN)	0003000
4.	INTEGER A (*,N),IR(*),IZ(N)	0004000
5.	70 FORMAT(1H5,7X,25I5)	0005000
6.	80 FORMAT(1H5,10,3X,25I5)	0006000
7.	MM=140(25,MM)	0007000
8.	DP 10 IB1,MP	0008000
9.	10 IR(I)=R(I)+1,E=5+C,1	1009000
10.	DD 30 JB1,NN	0010000
11.	20 IZ(J)=Z(J)+1,E=2+C,1	0011000
12.	PRINT 70	0012000
13.	PRINT 70,(IR(I),IB1,MP)	0013000
14.	PRINT 70	0014000
15.	DD 30 JJ=1,NN	0015000
16.	JBN=N1=JJ	0016000
17.	30 PRINT 80,IZ(J),(A(I,J),IB1,MP)	0017000
18.	RETURN	0018000
19.	END	0019000

END

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